

Practica 3

2.

$$U = \frac{a_0 + i \vec{\sigma} \cdot \vec{a}}{a_0 - i \vec{\sigma} \cdot \vec{a}}$$

$$\vec{\sigma} \cdot \vec{a} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} a_1 + \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} a_2 + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} a_3$$

a)

$$U^\dagger U = 1$$

$$U^\dagger U = \frac{(a_0 + i \vec{\sigma} \cdot \vec{a})^\dagger (a_0 + i \vec{\sigma} \cdot \vec{a})}{(a_0 - i \vec{\sigma} \cdot \vec{a}) (a_0 - i \vec{\sigma} \cdot \vec{a})}$$

$$U^\dagger U = 1$$

⇒ U es unitaria

$$U = \frac{a_0 + i \begin{pmatrix} a_3 & a_1 - i a_2 \\ a_1 + i a_2 & -a_3 \end{pmatrix}}{a_0 - i \begin{pmatrix} a_3 & a_1 - i a_2 \\ a_1 + i a_2 & -a_3 \end{pmatrix}}$$

$$U^\dagger = \frac{a_0 - i \begin{pmatrix} a_3 & a_1 - i a_2 \\ a_1 + i a_2 & -a_3 \end{pmatrix}}{a_0 + i \begin{pmatrix} a_3 & a_1 - i a_2 \\ a_1 + i a_2 & -a_3 \end{pmatrix}}$$

$$\rightarrow (\vec{\sigma} \cdot \vec{a})^\dagger = \vec{\sigma} \cdot \vec{a}$$

$$U = \frac{\begin{pmatrix} a_0 & 0 \\ 0 & a_0 \end{pmatrix} + i \begin{pmatrix} a_3 & a_1 - i a_2 \\ a_1 + i a_2 & -a_3 \end{pmatrix}}{\begin{pmatrix} a_0 & 0 \\ 0 & a_0 \end{pmatrix} - i \begin{pmatrix} a_3 & a_1 - i a_2 \\ a_1 + i a_2 & -a_3 \end{pmatrix}}$$

$$\rightarrow \det = \left\{ (a_0 + i a_3)(a_0 - i a_3) - (a_2 + i a_1)(-a_2 + i a_1) \right\}$$

$$a_0^2 + i a_0 a_3 - i a_0 a_3 - a_3^2 - (-a_2^2 - i a_1 a_2 - a_2^2 + i a_1 a_2)$$

$$\det = \{ (a_0 - i a_3)(a_0 + i a_3) - (-a_2 - i a_1)(+a_2 - i a_1) \}$$

$$= \{ a_0^2 - i a_0 a_3 + i a_0 a_3 - a_3^2 - (-a_2^2 - i a_1 a_2 + i a_1 a_2 - a_1^2) \} = a_0^2 + a_3^2 + a_2^2 + a_1^2$$

$$U = \frac{a_0 + i \vec{\sigma} \cdot \vec{a}}{a_0 - i \vec{\sigma} \cdot \vec{a}} \cdot \frac{(a_0 + i \vec{\sigma} \cdot \vec{a})}{(a_0 + i \vec{\sigma} \cdot \vec{a})} = \frac{a_0^2 + 2a_0 i \vec{\sigma} \cdot \vec{a} - (\vec{\sigma} \cdot \vec{a})^2}{a_0^2 + (\vec{\sigma} \cdot \vec{a})^2}$$

$$= \frac{(a_0 + i \vec{\sigma} \cdot \vec{a})^2}{a_0^2 + |\vec{a}|^2} = \frac{a_0^2 - |\vec{a}|^2 + 2a_0 i (\vec{\sigma} \cdot \vec{a})}{a_0^2 + a_1^2 + a_2^2 + a_3^2}$$

$$\frac{\det(a_0 + i \vec{\sigma} \cdot \vec{a})}{\det(a_0 - i \vec{\sigma} \cdot \vec{a})} = 1$$

⇒ U tiene det = 1

$$U = \frac{(a_0^2 - |\vec{a}|^2) \mathbb{1} + i (\vec{\sigma} \cdot \vec{a}) 2a_0}{(a_0^2 + |\vec{a}|^2)}$$

U es unimodular

$$U = \frac{a_0^2 - a_1^2 - a_2^2 - a_3^2}{a_0^2 + a_1^2 + a_2^2 + a_3^2} \mathbb{1} + i \frac{(\vec{\sigma} \cdot \vec{a}) 2a_0}{|\vec{a}| a_0^2 + |\vec{a}|^2}$$

$$\frac{\varphi}{2} = \arcsen \left(\frac{2a_0 |\vec{a}|}{a_0^2 + |\vec{a}|^2} \right)$$

$$\left(\frac{a_0^2 - |\vec{a}|^2}{a_0^2 + |\vec{a}|^2} \right)^2 + \left(\frac{2a_0 |\vec{a}|}{a_0^2 + |\vec{a}|^2} \right)^2 = 1$$

$$(a_0^2 - |\vec{a}|^2)^2 + 4a_0^2 |\vec{a}|^2 = (a_0^2 + |\vec{a}|^2)^2$$

$$\hat{a} = \hat{n} = \frac{a_1}{|\vec{a}|} \hat{x} + \frac{a_2}{|\vec{a}|} \hat{y} + \frac{a_3}{|\vec{a}|} \hat{z}$$

$$\frac{\varphi}{2} = \arcsen \left(\frac{2a_0 |\vec{a}|}{a_0^2 + |\vec{a}|^2} \right) = \arccos \left(\frac{a_0^2 - |\vec{a}|^2}{a_0^2 + |\vec{a}|^2} \right)$$

Esta matriz U representa una rotación $e^{-i \vec{\sigma} \cdot \vec{a} \frac{\varphi}{2}}$

4.

Partícula de spin 1

$$A \equiv S_z (S_x + \hbar) (S_x - \hbar)$$

$$\left. \begin{aligned} (S_x - \hbar)|+\rangle &= 0 \\ (S_x + \hbar)(2\hbar)|-\rangle &= 0 \\ S_z(\hbar)(\hbar)|0\rangle &= 0 \end{aligned} \right\}$$

Todos los elementos son nulos

$$(A)_{jk} = 0$$

$$|+\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \quad |-\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \quad |0\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$$

$$S_z \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = \hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = 0$$

$$J_x \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix} = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix}$$

$$\hat{B} \equiv S_x (S_x + \hbar) (S_x - \hbar)$$

$$\hat{B}|0\rangle$$

$$(S_x + \hbar) \left(\frac{\hbar}{2}|-\rangle - \hbar|+\rangle \right)$$

$$(S_x + \hbar) \left(\frac{\hbar}{2}|+\rangle - \hbar|-\rangle \right)$$

$$-\frac{\hbar}{\sqrt{2}} \frac{\hbar}{2} (|-\rangle + |+\rangle)$$

$$S_x \left(\frac{\hbar^2}{2}|-\rangle + \frac{\hbar^2}{4}|+\rangle + \frac{\hbar^2}{2}|-\rangle - \hbar^2|+\rangle \right) =$$

$$S_x \left(\frac{\hbar^2}{4}|-\rangle + \frac{\hbar^2}{2}|+\rangle - \hbar^2|-\rangle - \frac{\hbar^2}{2}|+\rangle \right)$$

$$B = \begin{pmatrix} \frac{\hbar^3}{\sqrt{2}} & 0 & 0 \\ 0 & \frac{3\hbar^3}{8} & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\left(\frac{\hbar^3}{4}|+\rangle + \frac{\hbar^3}{8}|-\rangle + \frac{\hbar^3}{4}|+\rangle - \frac{\hbar^3}{2}|-\rangle \right)$$

$$B = \frac{3\hbar^3}{8} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

5.

$$H = \frac{1}{2} \left(\frac{K_1^2}{I_1} + \frac{K_2^2}{I_2} + \frac{K_3^2}{I_3} \right)$$

Hamiltoniano de un cuerpo rígido

$$\frac{dK_1}{dt} = \frac{1}{i\hbar} [K_1, H] = \frac{1}{i\hbar} \left([K_1, \frac{K_2^2}{2I_1}] + [K_1, \frac{K_3^2}{2I_3}] \right)$$

$$= \frac{1}{i\hbar} \left(\frac{1}{2I_2} (K_2 [K_1, K_2] + [K_1, K_2] K_2) + \frac{1}{2I_3} (K_3 [K_1, K_3] + [K_1, K_3] K_3) \right)$$

$$\frac{dK_1}{dt} = \frac{1}{i\hbar} \left(\frac{1}{2I_2} [K_2, i\hbar K_3 + i\hbar K_3 K_2] + \frac{1}{2I_3} [K_3, i\hbar K_2 + -i\hbar K_2 K_3] \right)$$

$$\frac{dK_1}{dt} = \frac{1}{2I_2} (K_2 K_3 + K_3 K_2) - \frac{1}{2I_3} (K_3 K_2 + K_2 K_3)$$

en forma idem deberíamos llegar a:

Clásicamente

$$\frac{dK_1}{dt} = \frac{K_2 K_3}{I_2} - \frac{K_2 K_3}{I_3} = K_2 K_3 \left(\frac{I_3 - I_2}{I_2 I_3} \right)$$

$$\frac{dK_1}{dt} = \frac{K_2 K_3 (I_3 - I_2)}{I_2 I_3}$$

$$\frac{dK_2}{dt} = \frac{1}{i\hbar} \left([K_2, \frac{K_1^2}{2I_1}] + [K_2, \frac{K_3^2}{2I_3}] \right)$$

$$= \frac{1}{i\hbar} \left(\frac{1}{2I_1} (K_1 [K_2, K_1] + [K_2, K_1] K_1) + \frac{1}{2I_3} (K_3 [K_2, K_3] + [K_2, K_3] K_3) \right)$$

$$\frac{dK_2}{dt} = \frac{K_1 i\hbar K_2 + i\hbar K_2 K_1}{2I_1} - \frac{K_3 i\hbar K_1 - i\hbar K_1 K_3}{2I_3}$$

asimétrico

clásicamente

$$\frac{dK_2}{dt} = \frac{K_1 K_3 (I_2 - I_1)}{I_1 I_2}$$

$$\frac{dK_3}{dt} = \frac{K_1 K_2 (I_3 - I_1)}{I_1 I_3}$$

7.

$$|\alpha\rangle = \sum_a |a'\rangle \langle a'|\alpha\rangle \quad \text{base } \{|+\rangle, |-\rangle\}$$

$$|\alpha\rangle = |+\rangle \langle +|\alpha\rangle + |-\rangle \langle -|\alpha\rangle$$

$$D(\hat{z}, \varphi) = e^{-i \frac{S_z}{\hbar} \varphi} \rightarrow$$

$$|\alpha\rangle_R = e^{-i S_z \varphi / \hbar} |\alpha\rangle$$

$$S_z = \frac{\hbar}{2} (|+\rangle \langle +| - |-\rangle \langle -|)$$

$$S_x |+\rangle = \frac{\hbar}{2} |-\rangle$$

$$S_x |-\rangle = -\frac{\hbar}{2} |+\rangle$$

$$= \sum_{n=0}^{\infty} \left(\frac{i\varphi}{\hbar}\right)^n \frac{S_z^n}{n!} (|+\rangle \langle +|\alpha\rangle + |-\rangle \langle -|\alpha\rangle)$$

$$= \sum_{n=0}^{\infty} \left(\frac{i\varphi}{\hbar}\right)^n \left(\frac{\hbar}{2}\right)^n (|+\rangle \langle +|\alpha\rangle + |-\rangle \langle -|\alpha\rangle) + \sum_{n=0}^{\infty} \left(\frac{i\varphi}{\hbar}\right)^n \left(\frac{\hbar}{2}\right)^n (|-\rangle \langle -|\alpha\rangle)$$

$$|\alpha\rangle_R = e^{-i\varphi/2} |+\rangle \langle +|\alpha\rangle + e^{i\varphi/2} |-\rangle \langle -|\alpha\rangle \quad \leftarrow \text{estado arbitrario rotado}$$

a).

$$\langle \alpha | S_x | \alpha \rangle_R = \langle S_x \rangle_R$$

$$= \left(\langle \alpha | + \rangle \langle + | e^{i\varphi/2} + \langle \alpha | - \rangle \langle - | e^{-i\varphi/2} \right) \frac{\hbar}{2} (|+\rangle \langle +| + |-\rangle \langle -|) \left(e^{-i\varphi/2} |+\rangle \langle +|\alpha\rangle + e^{i\varphi/2} |-\rangle \langle -|\alpha\rangle \right)$$

$$= \left(\dots \right) \frac{\hbar}{2} (|-\rangle \langle +|\alpha\rangle e^{-i\varphi/2} + |+\rangle \langle -|\alpha\rangle e^{i\varphi/2})$$

$$= \frac{\hbar}{2} (\langle \alpha | + \rangle \langle -|\alpha\rangle e^{i\varphi} + \langle \alpha | - \rangle \langle +|\alpha\rangle e^{-i\varphi})$$

$$= \frac{\hbar}{2} (\langle \alpha | + \rangle \langle -|\alpha\rangle + \langle \alpha | - \rangle \langle +|\alpha\rangle) \cos \varphi + i (\langle \alpha | + \rangle \langle -|\alpha\rangle - \langle \alpha | - \rangle \langle +|\alpha\rangle) \sin \varphi$$

$$\langle S_x \rangle_R = \frac{\hbar}{2} \left\{ \langle \alpha | (|+\rangle \langle +| + |-\rangle \langle -|) | \alpha \rangle \right\} \cos \varphi + i \sin \varphi \left\{ \langle \alpha | (|+\rangle \langle +| - |-\rangle \langle -|) | \alpha \rangle \right\}$$

$$\langle S_x \rangle_R = \langle \alpha | S_x | \alpha \rangle \cdot \cos \varphi - \langle \alpha | S_y | \alpha \rangle \cdot \sin \varphi$$

$$\boxed{\langle S_x \rangle_R = \langle S_x \rangle \cdot \cos \varphi - \langle S_y \rangle \sin \varphi}$$

Los valores de expectación transforman como vectores

b)

$$D(i, 2\pi) |\alpha\rangle = -|+\rangle \langle +|\alpha\rangle - |-\rangle \langle -|\alpha\rangle = -|\alpha\rangle$$

$$\boxed{|\alpha\rangle_R = -|\alpha\rangle}$$

$$|\alpha\rangle_R = e^{i\pi} |\alpha\rangle$$

factor de fase

No se obtiene el mismo resultado sino el opuesto. Puede verse que si rotamos en 4π sí obtenemos el mismo resultado. (hay que dar dos vueltas)

8.

$$a) \quad (\vec{\sigma} \cdot \vec{a})(\vec{\sigma} \cdot \vec{b}) = \vec{a} \cdot \vec{b} + i \vec{\sigma} (\vec{a} \times \vec{b})$$

$$(\sigma_1 a_1 + \sigma_2 a_2 + \sigma_3 a_3)(\sigma_1 b_1 + \sigma_2 b_2 + \sigma_3 b_3) =$$

$$\sum_j \sigma_j a_j \sum_k \sigma_k b_k = \sum_j a_j b_j + \sum_k \overbrace{i \sigma_k \epsilon_{ijk} a_j b_k}^{(a \times b)_i}$$

$$\frac{1}{2} \underbrace{2i \sigma_k \epsilon_{ijk}}_{[\sigma_j, \sigma_k]} = \frac{1}{2} (2i \sigma_k \epsilon_{jki})$$

$$\sum_{j,k} a_j b_k \delta_{jk} + \sum_j \sum_k \frac{1}{2} [\sigma_j, \sigma_k] a_j b_k$$

$$= \sum_{j,k} a_j b_k \left(\delta_{jk} + \frac{1}{2} [\sigma_j, \sigma_k] \right)$$

$$\sum_{j,k} \sigma_j a_j \sigma_k b_k = \sum_{j,k} a_j b_k \left(\frac{1}{2} \{ \sigma_j, \sigma_k \} + \frac{1}{2} [\sigma_j, \sigma_k] \right)$$

$$\sum_{j,k} \sigma_j \sigma_k a_j b_k = \sum_{j,k} a_j b_k \left(\frac{\sigma_j \sigma_k + \sigma_j \sigma_k + \sigma_j \sigma_k - \sigma_j \sigma_k}{2} \right)$$

$$= \sum_{j,k} a_j b_k \cdot \sigma_j \sigma_k \rightarrow \boxed{\text{vale la igualdad}}$$

b) $\mathcal{D}(\hat{n}, \phi) = e^{-i \frac{\vec{\sigma} \cdot \hat{n} \phi}{\hbar}}$

$$\vec{\sigma} \cdot \hat{n} = \frac{\hbar}{2} (\sigma_1 n_1 + \sigma_2 n_2 + \sigma_3 n_3)$$

$$\text{con } \begin{cases} \vec{\sigma} = (\sigma_1, \sigma_2, \sigma_3) \\ \hat{n} = (n_1, n_2, n_3) \end{cases}$$

$$\vec{\sigma} = \frac{\hbar}{2} \vec{\sigma}$$

$$\mathcal{D}(\hat{n}, \phi) = e^{-i \frac{\vec{\sigma} \cdot \hat{n} \phi}{2}}$$

$$\mathcal{D}(\hat{n}, \phi) = \sum_{n=0}^{\infty} \frac{1}{n!} \left(-i \frac{\vec{\sigma} \cdot \hat{n} \phi}{2} \right)^n = \sum_{n=0}^{\infty} \frac{1}{(2n)!} \frac{(-i)^{2n} (\vec{\sigma} \cdot \hat{n})^{2n} \phi^{2n}}{2^{2n}} + \sum_{n=0}^{\infty} \frac{1}{(2n+1)!} \frac{(-i)^{2n+1} (\vec{\sigma} \cdot \hat{n})^{2n+1} \phi^{2n+1}}{2^{2n+1}}$$

$$\mathcal{D}(\hat{n}, \phi) = \sum_{n=0}^{\infty} \frac{(-1)^n (\phi/2)^{2n}}{(2n)!} + (-i \vec{\sigma} \cdot \hat{n}) \sum_{n=0}^{\infty} \frac{(-1)^n (\phi/2)^{2n+1}}{(2n+1)!}$$

$$(\vec{\sigma} \cdot \hat{n})^2 = |\hat{n}|^2 = 1$$

$$(\vec{\sigma} \cdot \hat{n})^4 = [(\vec{\sigma} \cdot \hat{n})^2]^2 = 1^2 = 1$$

$$(\vec{\sigma} \cdot \hat{n})^{2n} = [(\vec{\sigma} \cdot \hat{n})^2]^n = 1^n = 1 \quad \forall n \in \mathbb{N}$$

$$(\vec{\sigma} \cdot \hat{n})^{2n+1} = (\vec{\sigma} \cdot \hat{n})^{2n} \vec{\sigma} \cdot \hat{n} = \vec{\sigma} \cdot \hat{n} \quad \forall n \in \mathbb{N}$$

$$(-i)^{2n} = [(-i)^2]^n = (-1)^n \Rightarrow$$

$$\mathcal{D}(\hat{n}, \phi) = \sum_{n=0}^{\infty} \frac{(-1)^n (\phi/2)^{2n}}{(2n)!}$$

$$-i \vec{\sigma} \cdot \hat{n} \sum_{n=0}^{\infty} \frac{(-1)^n (\phi/2)^{2n+1}}{(2n+1)!}$$

usando la parte a)

$$\mathcal{D}(\hat{n}, \phi) = \mathbb{1} \cos\left(\frac{\phi}{2}\right) - i \vec{\sigma} \cdot \hat{n} \text{sen}\left(\frac{\phi}{2}\right)$$

donde $\mathbb{1}$ es la matriz identidad

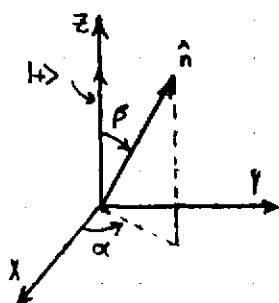
c)

$$D(\hat{n}, \phi) = \begin{pmatrix} \cos(\phi/2) & 0 \\ 0 & \cos(\phi/2) \end{pmatrix} - i \begin{pmatrix} n_3 & n_1 - i n_2 \\ n_1 + i n_2 & -n_3 \end{pmatrix} \sin(\frac{\phi}{2})$$

$$D(\hat{n}, \phi) = \begin{pmatrix} \cos(\phi/2) - i n_3 \sin(\phi/2) & (-i n_1 - n_2) \sin(\phi/2) \\ (-i n_1 + n_2) \sin(\phi/2) & \cos(\phi/2) + i n_3 \sin(\phi/2) \end{pmatrix}$$

d) $\hat{n} = \hat{x} \cos \alpha \sin \beta + \hat{y} \sin \alpha \sin \beta + \hat{z} \cos \beta$

$|\hat{S} \cdot \hat{n}; +\rangle$



$$\hat{S} \cdot \hat{n} = S_x \cos \alpha \sin \beta + S_y \sin \alpha \sin \beta + S_z \cos \beta$$

1. Hay que rotar β en torno a \hat{y}
2. Hay que rotar α en torno a \hat{z}

$$D(\hat{y}, \beta) = \mathbb{1} \cos(\beta/2) - i \sigma_y \sin(\beta/2)$$

$$D(\hat{z}, \alpha) = \mathbb{1} \cos(\alpha/2) - i \sigma_z \sin(\alpha/2)$$

$$|\hat{S} \cdot \hat{n}; +\rangle = D(\hat{z}, \alpha) \cdot D(\hat{y}, \beta) |+\rangle = \begin{pmatrix} \cos(\alpha/2) - i \sin(\alpha/2) & 0 \\ 0 & \cos(\alpha/2) + i \sin(\alpha/2) \end{pmatrix} \begin{pmatrix} \cos(\beta/2) & 0 \\ 0 & \cos(\beta/2) \end{pmatrix} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \sin(\beta/2) \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$|\hat{S} \cdot \hat{n}; +\rangle = \begin{pmatrix} e^{-i\alpha/2} & 0 \\ 0 & e^{i\alpha/2} \end{pmatrix} \begin{pmatrix} \cos(\beta/2) & -\sin(\beta/2) \\ \sin(\beta/2) & \cos(\beta/2) \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$\begin{pmatrix} \cos(\beta/2) \\ \sin(\beta/2) \end{pmatrix} \Rightarrow |\hat{S} \cdot \hat{n}; +\rangle = e^{-i\alpha/2} \cos(\frac{\beta}{2}) |+\rangle + e^{i\alpha/2} \sin(\frac{\beta}{2}) |-\rangle$$

9.

$$D^{1/2}(\alpha, \beta, \gamma) = D(\hat{z}, \alpha) \cdot D(\hat{y}, \beta) \cdot D(\hat{z}, \gamma)$$

a)

$$D^{1/2}(\alpha, \beta, \gamma) = e^{-\frac{i \hat{J}_z \alpha}{\hbar}} e^{-\frac{i \hat{J}_y \beta}{\hbar}} e^{-\frac{i \hat{J}_z \gamma}{\hbar}}$$

Pero \rightarrow

$$\frac{\hat{J} \cdot \hat{n}}{\hbar} = \frac{J_n}{\hbar} = \frac{S_n}{\hbar} = \frac{\hbar \sigma_n}{2\hbar}$$

$$D^{1/2}(\alpha, \beta, \gamma) = e^{-\frac{i \sigma_z \alpha}{2}} e^{-\frac{i \sigma_y \beta}{2}} e^{-\frac{i \sigma_z \gamma}{2}}$$

$$\mathbb{1} \cos(\frac{\alpha}{2}) - i \sigma_z \sin(\frac{\alpha}{2}) \rightarrow \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\mathbb{1} \cos(\frac{\beta}{2}) - i \sigma_y \sin(\frac{\beta}{2}) \rightarrow \begin{pmatrix} \cos(\beta/2) & -i \sin(\beta/2) \\ i \sin(\beta/2) & \cos(\beta/2) \end{pmatrix}$$

$$\mathbb{1} \cos(\frac{\gamma}{2}) - i \sigma_z \sin(\frac{\gamma}{2}) \rightarrow \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

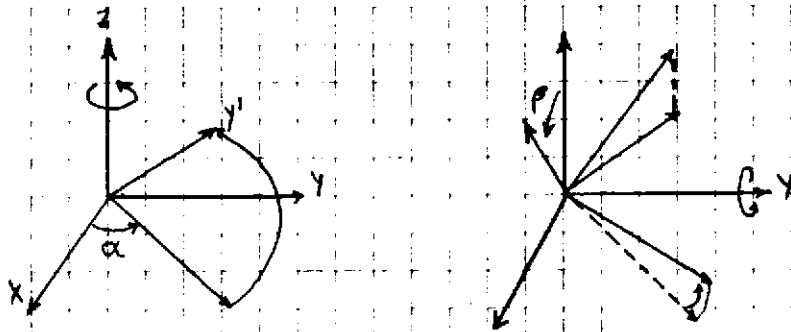
$$\begin{pmatrix} \cos(\frac{\alpha}{2}) - i \sin(\frac{\alpha}{2}) & 0 \\ 0 & \cos(\frac{\alpha}{2}) + i \sin(\frac{\alpha}{2}) \end{pmatrix} \begin{pmatrix} \cos(\frac{\beta}{2}) & -\sin(\frac{\beta}{2}) \\ \sin(\frac{\beta}{2}) & \cos(\frac{\beta}{2}) \end{pmatrix} \begin{pmatrix} \cos(\frac{\gamma}{2}) - i \sin(\frac{\gamma}{2}) & 0 \\ 0 & \cos(\frac{\gamma}{2}) + i \sin(\frac{\gamma}{2}) \end{pmatrix}$$

$$\begin{pmatrix} e^{-i\alpha/2} & 0 \\ 0 & e^{i\alpha/2} \end{pmatrix} \begin{pmatrix} \cos(\beta/2) & -\sin(\beta/2) \\ \sin(\beta/2) & \cos(\beta/2) \end{pmatrix} \begin{pmatrix} e^{-i\gamma/2} & 0 \\ 0 & e^{i\gamma/2} \end{pmatrix}$$

$$\begin{pmatrix} e^{i\frac{\alpha\beta}{2}} \cos(\beta/2) & -e^{i\alpha/2} \sin(\beta/2) \\ e^{i\alpha/2} \sin(\beta/2) & e^{i\alpha\beta/2} \cos(\beta/2) \end{pmatrix} \begin{pmatrix} e^{-i\gamma/2} & 0 \\ 0 & e^{i\gamma/2} \end{pmatrix}$$

$$\mathcal{D}_{(\alpha, \beta, \gamma)}^{1/2} = \begin{pmatrix} e^{-\frac{i\alpha}{2} - \frac{i\gamma}{2}} \cos(\frac{\beta}{2}) & -e^{-\frac{i\alpha}{2} + \frac{i\gamma}{2}} \sin(\frac{\beta}{2}) \\ e^{\frac{i\alpha}{2} - \frac{i\gamma}{2}} \sin(\frac{\beta}{2}) & e^{\frac{i\alpha}{2} + \frac{i\gamma}{2}} \cos(\frac{\beta}{2}) \end{pmatrix}$$

b)



Usando lo hecho en ejercicio 8c tenemos:

$$\mathcal{D}_{(\hat{n}, \theta)}^{1/2} = \begin{pmatrix} \cos(\theta/2) - i n_3 \sin(\theta/2) & (-i n_1 - n_2) \sin(\theta/2) \\ (-i n_1 + n_2) \sin(\theta/2) & \cos(\theta/2) + i n_3 \sin(\theta/2) \end{pmatrix}$$

Le podemos sacar fase global fuera \rightarrow

$$\mathcal{D}^{1/2} = e^{-i\alpha/2} \begin{pmatrix} 1 & 0 \\ 0 & e^{i\alpha} \end{pmatrix} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \begin{pmatrix} 1 \\ e^{i\gamma} \end{pmatrix} e^{-i\gamma/2} \mathcal{D}^{1/2} = e^{-i\alpha/2 - i\gamma/2} \begin{pmatrix} \cos(\frac{\beta}{2}) & -e^{i\gamma} \sin(\frac{\beta}{2}) \\ e^{i\alpha} \sin(\frac{\beta}{2}) & e^{i\alpha\gamma} \cos(\frac{\beta}{2}) \end{pmatrix}$$

$$\left(\cos^2(\theta/2) + n_3^2 \sin^2(\theta/2) \right) + \left(\sin^2(\theta/2) \right) (n_1^2 + n_2^2) = n_1^2 + n_2^2 + n_3^2 = 1$$

$$\begin{cases} \cos(\frac{\theta}{2}) = \cos(\frac{\beta}{2}) \cos(\frac{\alpha+\gamma}{2}) \\ n_3 \sin(\frac{\theta}{2}) = \cos(\frac{\beta}{2}) \sin(\frac{\alpha+\gamma}{2}) \\ n_2 \sin(\frac{\theta}{2}) = \sin(\frac{\beta}{2}) \cos(\frac{\alpha+\gamma}{2}) \\ n_1 \sin(\frac{\theta}{2}) = \sin(\frac{\beta}{2}) \sin(\frac{\alpha+\gamma}{2}) \end{cases}$$

datos α, β, γ

incógnitas θ, n_1, n_2, n_3

$$1 = \cos^2(\frac{\beta}{2}) \cos^2(\frac{\alpha+\gamma}{2}) + \frac{1}{n_3^2} \cos^2(\frac{\beta}{2}) \sin^2(\frac{\alpha+\gamma}{2})$$

$$\frac{1}{n_3^2} = \frac{1 - \cos^2(\beta/2) \cdot \cos^2(\alpha+\gamma/2)}{\cos^2(\beta/2) \cdot \sin^2(\alpha+\gamma/2)}$$

$$1 = \cos^2\left(\frac{\beta}{2}\right) \cdot \cos^2\left(\frac{\alpha+\gamma}{2}\right) + \frac{1}{n_2^2} \cdot \sin^2\left(\frac{\beta}{2}\right) \cdot \cos^2\left(\frac{\alpha-\gamma}{2}\right)$$

$$\frac{1}{n_2^2} = \frac{1 - \cos^2(\beta/2) \cdot \cos^2(\alpha+\gamma/2)}{\sin^2(\beta/2) \cdot \cos^2(\alpha-\gamma/2)}$$

$$1 = \cos^2\left(\frac{\beta}{2}\right) \cdot \cos^2\left(\frac{\alpha+\gamma}{2}\right) + \frac{1}{n_1^2} \cdot \sin^2\left(\frac{\beta}{2}\right) \cdot \cos^2\left(\frac{\alpha-\gamma}{2}\right)$$

$$\frac{1}{n_1^2} = \frac{1 - \cos^2(\beta/2) \cdot \cos^2(\alpha+\gamma/2)}{\sin^2(\beta/2) \cdot \cos^2(\alpha-\gamma/2)}$$

$$n_1^2 = n_2^2 \rightarrow n_1 = n_2 \rightarrow n_3^2 = 2n_1^2 \rightarrow n_3 = \sqrt{2} n_1$$

$$\frac{\sin(\theta)}{2} = \sin\left(\frac{\theta}{2}\right) \cos\left(\frac{\theta}{2}\right) = \cos\left(\frac{\beta}{2}\right) \cdot \cos\left(\frac{\alpha+\gamma}{2}\right) \cdot \frac{1}{n_2} \cdot \sin\left(\frac{\beta}{2}\right) \cdot \cos\left(\frac{\alpha-\gamma}{2}\right)$$

$$\frac{\sin(\theta)}{2} = \frac{\sin(\beta)}{2} \cdot \frac{\sqrt{1 - \cos^2(\beta/2) \cdot \cos^2(\alpha+\gamma/2)} \cdot \cos(\alpha-\gamma/2) \cdot \cos(\alpha+\gamma/2)}{\sin(\beta/2) \cdot \cos(\alpha-\gamma/2)}$$

$$\sqrt{\sin^2(\beta/2) + \cos^2(\beta/2) [1 - \cos^2(\alpha+\gamma/2)]}$$

$$\sin(\theta) = \frac{\sin(\beta)}{\sin(\beta/2)} \cdot \sqrt{\sin^2(\beta/2) + \cos^2(\beta/2) \cdot \sin^2(\alpha+\gamma/2)}$$

$$\cos(\theta) = \cos\left(\frac{\beta}{2}\right) \cdot \cos\left(\frac{\alpha+\gamma}{2}\right)$$

$$\sin\left(\frac{\beta}{2} + \frac{\beta}{2}\right) = 2 \sin\left(\frac{\beta}{2}\right) \cdot \cos\left(\frac{\beta}{2}\right)$$

$$\sin \theta = 2 \cdot \cos\left(\frac{\beta}{2}\right) \cdot \sqrt{\sin^2\left(\frac{\beta}{2}\right) + \cos^2\left(\frac{\beta}{2}\right) \cdot \sin^2\left(\frac{\alpha+\gamma}{2}\right)}$$

$$\cos \theta = \cos\left(\frac{\beta}{2}\right) \cdot \cos\left(\frac{\alpha+\gamma}{2}\right)$$

$$\sin \theta = \frac{2 \cdot \cos \theta \cdot \sqrt{1 - \cos^2(\beta/2) \cdot \cos^2(\alpha+\gamma/2)}}{\cos(\alpha+\gamma/2)}$$

$$\tan \theta = 2 \cdot \sqrt{\frac{1}{\cos^2(\alpha+\gamma/2)} - \cos^2(\beta/2)}$$

$$\theta = \arctan \left\{ 2 \cdot \left(\cos^{-2}\left[\frac{\alpha+\gamma}{2}\right] - \cos^2\left[\frac{\beta}{2}\right] \right)^{1/2} \right\}$$

$$J_+ |j, m\rangle = \langle + | j, m+1 \rangle \rightarrow \langle j, m | (J_+)^{\dagger} = \langle j, m+1 | C_+^*$$

$$\langle j, m | \underbrace{J_+}_{(J_+)^{\dagger}} | j, m \rangle = C_+^* C_+ \langle j, m+1 | j, m+1 \rangle$$

$$\langle j, m | J_+ | j, m \rangle = \hbar^2 (j(j+1) - m(m+1)) \langle j, m | j, m \rangle$$

$$\Rightarrow |C_+|^2 = \hbar^2 (j(j+1) - m(m+1))$$

Procediendo de modo idem con $J_- J_- = (J_-)^{\dagger} J_-$

$$J_{\pm} |j, m\rangle = \hbar (j(j+1) - m(m \pm 1))^{\frac{1}{2}} |j, m \pm 1\rangle$$

$$J_+ |j, j\rangle = \hbar (j(j+1) - j(j+1))^{\frac{1}{2}} |j, j+1\rangle \rightarrow \boxed{J_+ |j, j\rangle = 0}$$

En forma idem será $\boxed{J_- |j, -j\rangle = 0}$

b)

$$J_z |j, m\rangle = m\hbar |j, m\rangle$$

$\langle \alpha | J_x | \alpha \rangle$, con $|\alpha\rangle$ autoestado de J_z
 $J_z |\alpha\rangle = \alpha |\alpha\rangle$

$$J_+ = J_x + iJ_y$$

$$J_- = J_x - iJ_y$$

$$\frac{J_+ + J_-}{2} = J_x$$

$$\frac{J_+ - J_-}{2i} = J_y$$

$$\langle J_x \rangle = \langle j, m | \frac{J_+ + J_-}{2} | j, m \rangle$$

$$\langle J_x \rangle = \frac{1}{2} \langle j, m | (|j, m+1\rangle + |j, m-1\rangle)$$

$$\boxed{\langle J_x \rangle = 0}$$

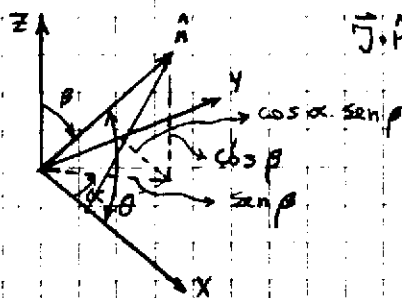
$$\langle J_y \rangle = \langle j, m | \frac{J_+ - J_-}{2i} | j, m \rangle$$

$$= \frac{1}{2i} \langle j, m | (|j, m+1\rangle - |j, m-1\rangle)$$

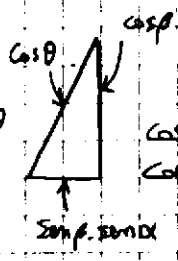
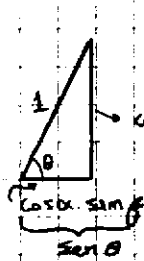
$$\boxed{\langle J_y \rangle = 0}$$

En $|\psi\rangle$ con $J_z |\psi\rangle = \hbar m |\psi\rangle$

$$\langle \psi | \vec{J} \cdot \hat{n} | \psi \rangle =$$



$$\vec{J} \cdot \hat{n} = J_x \cos \alpha \sin \beta + J_y \sin \alpha \sin \beta + J_z \cos \beta$$



$$\sin^2 \beta + \cos^2 \beta = 1$$

$$\cos^2 \theta + \cos^2 \alpha \sin^2 \beta = 1$$

$$\cos^2 \theta + \cos^2 \alpha (1 - \cos^2 \beta) = 1$$

$$\cos^2 \theta = \cos^2 \beta + \sin^2 \beta \sin^2 \alpha$$

$$\langle \vec{J} \cdot \hat{n} \rangle = \cos \alpha \sin \beta \langle \frac{J_+ + J_-}{2} \rangle + \sin \alpha \sin \beta \langle \frac{J_+ - J_-}{2i} \rangle + \langle J_z \rangle \cos \beta$$

$$\langle \vec{J} \cdot \hat{n} \rangle_\psi = \cos \alpha \cdot \underbrace{\sum m \langle J_x \rangle_\psi}_{=0} + \sin \alpha \cdot \underbrace{\sum m \langle J_y \rangle_\psi}_{=0} + \langle J_z \rangle \cos \beta$$

$$\begin{aligned} \langle \vec{J} \cdot \hat{n} \rangle_\psi &= \langle \psi | J_z | \psi \rangle \cdot \cos \beta \\ &= \hbar m \langle \psi | \psi \rangle \cdot \cos \beta \end{aligned}$$

$$\boxed{\langle \vec{J} \cdot \hat{n} \rangle_\psi = \hbar \cdot m \cdot \cos \beta}$$

Donde θ del enunciado debe ser ángulo entre \hat{z} y \hat{n} , no entre \hat{x} y \hat{n}

14.

base $\{|1,1\rangle, |1,0\rangle, |1,-1\rangle\}$

$$\begin{aligned} L^2 |j,m\rangle &= \hbar^2 j(j+1) |j,m\rangle \\ L_z |j,m\rangle &= m \hbar |j,m\rangle \end{aligned}$$

$$L^2 = \begin{pmatrix} 2\hbar^2 & 0 & 0 \\ 0 & 2\hbar^2 & 0 \\ 0 & 0 & 2\hbar^2 \end{pmatrix}$$

$$L_z = \begin{pmatrix} -\hbar \langle 1,-1|1,-1\rangle & 0 & \hbar \langle 1,-1|1,0\rangle \\ 0 & 0 & 0 \\ 0 & 0 & \hbar \langle 1,1|1,1\rangle \end{pmatrix} = \begin{pmatrix} -\hbar & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \hbar \end{pmatrix}$$

L^2, L_z son fáciles porque están en subbase de autoestados

$$\boxed{L^2 = 2\hbar^2 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} = 2\hbar^2 \mathbb{1}}$$

$$\boxed{L_z = \hbar \begin{pmatrix} -1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}}$$

$$L_x = \frac{L_+ + L_-}{2}$$

$$L_y = \frac{L_+ - L_-}{2i}$$

$$L_\pm |1, \pm 1\rangle = 0$$

los espines son máximos

$$L_y = \begin{pmatrix} 0 & +\frac{\sqrt{2}\hbar}{2i} & 0 \\ -\frac{\sqrt{2}\hbar}{2i} & \frac{\sqrt{2}\hbar}{2i} & \frac{\sqrt{2}\hbar}{2i} \\ 0 & -\frac{\sqrt{2}\hbar}{2i} & 0 \end{pmatrix}$$

$$\boxed{L_y = \frac{\hbar}{2} \begin{pmatrix} 0 & -\sqrt{2}i & 0 \\ \sqrt{2}i & 0 & -\sqrt{2}i \\ 0 & \sqrt{2}i & 0 \end{pmatrix}}$$

$$\langle 1,1 | \frac{L_+ - L_-}{2i} | 1,0 \rangle = \frac{\sqrt{2}\hbar}{2i}$$

$$\langle 1,0 | \frac{L_+ - L_-}{2i} | 1,-1 \rangle = \frac{\langle 1,0 |}{2i} (\hbar\sqrt{2} |1,0\rangle + \hbar\sqrt{2} |1,0\rangle)$$

$$\langle 1,0 | \frac{L_+ - L_-}{2i} | 1,1 \rangle = \frac{1}{2i} \langle 1,0 | (\sqrt{2}\hbar |1,1\rangle - \sqrt{2}\hbar |1,0\rangle)$$

$$\langle 1,-1 | \frac{L_+ - L_-}{2i} | 1,0 \rangle = \frac{1}{2i} \langle 1,-1 | (\sqrt{2}\hbar |1,1\rangle - \sqrt{2}\hbar |1,-1\rangle)$$

$$L_x = \begin{pmatrix} 0 & \frac{\hbar\sqrt{z}}{2} & 0 \\ \frac{\hbar\sqrt{z}}{2} & 0 & \frac{\hbar\sqrt{z}}{2} \\ 0 & \frac{\hbar\sqrt{z}}{2} & 0 \end{pmatrix}$$

$$L_x = \frac{\hbar}{2} \begin{pmatrix} 0 & \sqrt{z} & 0 \\ \sqrt{z} & 0 & \sqrt{z} \\ 0 & \sqrt{z} & 0 \end{pmatrix}$$

$L_x|1, \pm 1\rangle = 0$
Las espumas
son nulas.

$$\begin{aligned} \frac{1}{2} \langle 1, 0 | (L_+ + L_-) | 1, 1 \rangle &= \frac{1}{2} \langle 1, 0 | (\hbar\sqrt{2} | 1, 0 \rangle + 0) \\ \frac{1}{2} \langle 1, 0 | (L_+ + L_-) | 1, 0 \rangle &= \frac{1}{2} \langle 1, 0 | (| 1, 1 \rangle + | 1, -1 \rangle) \\ \frac{1}{2} \langle 1, 0 | (L_+ + L_-) | 1, -1 \rangle &= \frac{1}{2} \langle 1, 0 | (0 + \hbar\sqrt{2} | 1, 0 \rangle) \end{aligned}$$

$$\frac{1}{2} \langle 1, 1 | (L_+ + L_-) | 1, 0 \rangle = \frac{1}{2} \langle 1, 1 | (\hbar\sqrt{2} | 1, -1 \rangle)$$

$$[L_x, L_y] = L_x L_y - L_y L_x$$

$$\begin{aligned} &= \left(\frac{\hbar^2}{4}\right) \begin{pmatrix} 0 & \sqrt{z} & 0 \\ \sqrt{z} & 0 & \sqrt{z} \\ 0 & \sqrt{z} & 0 \end{pmatrix} \begin{pmatrix} 0 & -\sqrt{z}i & 0 \\ \sqrt{z}i & 0 & -\sqrt{z}i \\ 0 & \sqrt{z}i & 0 \end{pmatrix} - \frac{\hbar^2}{4} \begin{pmatrix} 0 & -\sqrt{z}i & 0 \\ \sqrt{z}i & 0 & -\sqrt{z}i \\ 0 & \sqrt{z}i & 0 \end{pmatrix} \begin{pmatrix} 0 & \sqrt{z} & 0 \\ \sqrt{z} & 0 & \sqrt{z} \\ 0 & \sqrt{z} & 0 \end{pmatrix} \\ &= \frac{\hbar^2}{4} \begin{pmatrix} z i & 0 & -z i \\ 0 & 0 & 0 \\ z i & 0 & -z i \end{pmatrix} - \frac{\hbar^2}{4} \begin{pmatrix} -z i & 0 & -z i \\ 0 & 0 & 0 \\ +z i & 0 & +z i \end{pmatrix} \end{aligned}$$

$$[L_x, L_y] = \frac{\hbar^2}{4} \begin{pmatrix} 4i & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -4i \end{pmatrix} = \hbar(\hbar i) \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix} = i\hbar L_z$$

esta OK

a) $\{|l, m_y\rangle\}$ base de L_z

$$L_y = \sqrt{z} \frac{\hbar}{2} \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix} = \frac{\hbar}{\sqrt{z}} \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix}$$

Hay que diagonalizar la matriz de L_y

$$\frac{\hbar}{\sqrt{z}} \begin{pmatrix} -\lambda & -1 & 0 \\ 1 & -\lambda & -1 \\ 0 & 1 & -\lambda \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix} = 0 \rightarrow ((-\lambda)^3 + 0 + 0) + (0 + \lambda + \lambda) =$$

autovalores de L_y

$$\lambda = 0 \rightarrow v_1 = \alpha \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix} \rightarrow \alpha = \frac{1}{\sqrt{2}}$$

$$\lambda = \sqrt{z}i$$

$$\begin{aligned} (-\sqrt{z}i)x - y &= 0 \\ y &= -\sqrt{z}ix \\ y - \sqrt{z}iz &= 0 \\ y &= \sqrt{z}iz \\ x &= -z \end{aligned}$$

$$v_2 = \begin{pmatrix} -z \\ \sqrt{z}iz \\ z \end{pmatrix} = \alpha \begin{pmatrix} 1 \\ \sqrt{z}i \\ 1 \end{pmatrix}$$

autovalores

$$\begin{aligned} \lambda &= 0 \\ \lambda &= +\sqrt{z}i \\ \lambda &= -\sqrt{z}i \end{aligned}$$

$$\begin{aligned} -\lambda^3 + 2\lambda &= 0 \\ (-\lambda^2 - 2)\lambda &= 0 \\ \lambda^2 &= -2 \\ \lambda &= \pm\sqrt{z}i \end{aligned}$$

$$\alpha^2 + \alpha^2 z + \alpha^2 = 1$$

$$4\alpha^2 = 1$$

$$\alpha^2 = \frac{1}{4} \rightarrow \alpha = \frac{1}{2}$$

$$\lambda = -\sqrt{z}i$$

$$\begin{aligned} \sqrt{z}ix - y &= 0 \\ y &= \sqrt{z}ix \\ y + \sqrt{z}iz &= 0 \\ y &= -\sqrt{z}iz \\ x &= -z \end{aligned}$$

$$v_3 = \begin{pmatrix} -z \\ \sqrt{z}iz \\ z \end{pmatrix} = \alpha \begin{pmatrix} -1 \\ -\sqrt{z}i \\ 1 \end{pmatrix} \rightarrow \alpha = \frac{1}{2}$$

AL definir la matriz de L^z según el arreglo:

$$\begin{pmatrix} \langle 1, -1 | & \langle 1, 0 | & \langle 1, 1 | & \\ \hline & & & | 1, -1 \rangle \\ & & & | 1, 0 \rangle \\ & & & | 1, 1 \rangle \\ \hline \end{pmatrix} \Rightarrow$$

hemos establecido

$$\begin{aligned} |1, -1\rangle &\equiv \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} && \text{asociado con } -\hbar \\ |1, 0\rangle &\equiv \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} && \text{asociado con } 0 \\ |1, 1\rangle &\equiv \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} && \text{asociado con } \hbar \end{aligned}$$

Base hallada del subespacio

$$B = \left\{ \underbrace{\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix}}_{v_1}, \underbrace{\frac{1}{2} \begin{pmatrix} -1 \\ \sqrt{2}i \\ 1 \end{pmatrix}}_{v_2}, \underbrace{\frac{1}{2} \begin{pmatrix} -1 \\ -\sqrt{2}i \\ 1 \end{pmatrix}}_{v_3} \right\}$$

$$v_1 = \frac{1}{\sqrt{2}} (|1, -1\rangle + |1, 1\rangle)$$

$$v_2 = \frac{1}{2} (-|1, -1\rangle + \sqrt{2}i |1, 0\rangle + |1, 1\rangle)$$

$$v_3 = \frac{1}{2} (-|1, -1\rangle - \sqrt{2}i |1, 0\rangle + |1, 1\rangle)$$

b) $|\psi\rangle = \frac{1}{\sqrt{2}} (|1, 1\rangle - |1, -1\rangle)$

Aquí confirmamos que los autovalores sean los que nos rechazamos según las asociaciones

$$L_y |v_1\rangle = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} \\ 0 \\ \frac{1}{\sqrt{2}} \end{pmatrix} = 0$$

$$L_y |v_2\rangle = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} \frac{-\sqrt{2}i}{2} \\ \frac{\sqrt{2}i}{2} \\ \frac{1}{2} \end{pmatrix}$$

$$\text{avalar } \hbar - \hbar \frac{1}{\sqrt{2}} \begin{pmatrix} -\frac{\sqrt{2}i}{2} \\ -1 \\ \frac{\sqrt{2}i}{2} \end{pmatrix}$$

$$L_y |v_3\rangle = +\hbar |v_3\rangle \quad \frac{1}{\sqrt{2}i} = \frac{\sqrt{2}i}{2}$$

Medimos $L_y \Rightarrow$ tener que proyectar $|\psi\rangle$ en la base de autoestados $\{v_1, v_2, v_3\} \Rightarrow$

$$\begin{aligned} P_{|v_1, 0\rangle} &= |\langle v_1 | \psi \rangle|^2 = \left| \frac{1}{\sqrt{2}} \cdot \frac{1}{\sqrt{2}} (\langle 1, -1 | + \langle 1, 1 |) (|1, 1\rangle - |1, -1\rangle) \right|^2 \\ &= \left| \frac{1}{2} (1 - 1) \right|^2 = \boxed{0} \quad \text{prob. de } 0 \end{aligned}$$

$$\begin{aligned} P_{|v_2, -\hbar\rangle} &= |\langle v_2 | \psi \rangle|^2 = \left| \frac{1}{2} \cdot \frac{1}{\sqrt{2}} (-\langle 1, -1 | - \langle 1, 0 | \sqrt{2}i + \langle 1, 1 |) (|1, 1\rangle - |1, -1\rangle) \right|^2 \\ &= \left| \frac{1}{2} \cdot \frac{1}{\sqrt{2}} \cdot 2 \right|^2 = \boxed{\frac{1}{2}} \quad \text{prob. de } -\hbar \end{aligned}$$

$$\begin{aligned} P_{|v_3, +\hbar\rangle} &= |\langle v_3 | \psi \rangle|^2 = \left| \frac{1}{2} \cdot \frac{1}{\sqrt{2}} (-\langle 1, -1 | + \langle 1, 0 | \sqrt{2}i + \langle 1, 1 |) (|1, 1\rangle - |1, -1\rangle) \right|^2 \\ &= \left| \frac{1}{2} \cdot \frac{1}{\sqrt{2}} \cdot (1 + 1) \right|^2 = \boxed{\frac{1}{2}} \quad \text{prob. de } +\hbar \end{aligned}$$

Medimos L_x

Ahora habría que diagonalizar L_x , para obtener su base de autoestados

$$\frac{\hbar}{\sqrt{2}} \begin{pmatrix} -\lambda & 1 & 0 \\ 1 & -\lambda & 1 \\ 0 & 1 & -\lambda \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix} = 0$$

$$((-\lambda)^3 + 0 + 0) - (0 - \lambda - \lambda) = 0$$

$$-\lambda^3 + 2\lambda = 0$$

$$\lambda^2 - 2 = 0$$

$$\lambda = 0$$

$$\lambda = \sqrt{2}$$

$$\lambda = -\sqrt{2}$$

* $\lambda = 0$; V_1
 $y = 0$
 $x + z = 0$

$$V_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix}$$

* $\lambda = \sqrt{2}$; V_2
 $-\sqrt{2}x + y = 0$
 $y = \sqrt{2}x$
 $y - \sqrt{2}z = 0$
 $y = \sqrt{2}z$

$$V_2 = \alpha \begin{pmatrix} 1 \\ \sqrt{2} \\ 1 \end{pmatrix}$$

$$\alpha^2(1 + 2 + 1) = 1$$

$$\alpha = \frac{1}{\sqrt{4}} = \frac{1}{2}$$

$$V_2 = \frac{1}{2} \begin{pmatrix} 1 \\ \sqrt{2} \\ 1 \end{pmatrix}$$

* $\lambda = -\sqrt{2}$; V_3
 $\sqrt{2}x + y = 0$
 $y = -\sqrt{2}x$
 $y + \sqrt{2}z = 0$
 $y = -\sqrt{2}z$

$$V_3 = \alpha \begin{pmatrix} 1 \\ -\sqrt{2} \\ 1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 \\ -\sqrt{2} \\ 1 \end{pmatrix}$$

$$V_0 = \frac{1}{\sqrt{2}} (|1,1\rangle - |1,-1\rangle); \quad V_2 = \frac{1}{2} (|1,1\rangle + \sqrt{2}|1,0\rangle + |1,-1\rangle)$$

$$V_3 = \frac{1}{2} (|1,1\rangle - \sqrt{2}|1,0\rangle + |1,-1\rangle)$$

$$\langle x | V_2 \rangle = \frac{\hbar}{2} \sqrt{2} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} \sqrt{2}/2 \\ \sqrt{2}/2 \\ \sqrt{2}/2 \end{pmatrix} = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} \sqrt{2}/2 \\ 1 \\ \sqrt{2}/2 \end{pmatrix}$$

$$\textcircled{+\hbar} \begin{pmatrix} \sqrt{2}/2 \\ \sqrt{2}/2 \end{pmatrix}$$

$$P_{|V_0, +\hbar\rangle} = \left| \frac{1}{2} (\langle 1,1| + \sqrt{2}\langle 1,0| + \langle 1,-1|) (|1,1\rangle - |1,-1\rangle) \frac{1}{\sqrt{2}} \right|^2 = \left| \frac{1}{2} \frac{1}{\sqrt{2}} (1 + (-1)) \right|^2 = \boxed{0}$$

no podemos medir $+\hbar$

$$P_{|V_2, -\hbar\rangle} = \left| \frac{1}{2} (\langle 1,-1| - \sqrt{2}\langle 1,0| + \langle 1,1|) (|1,1\rangle - |1,-1\rangle) \frac{1}{\sqrt{2}} \right|^2 = \left| \frac{1}{2} \frac{1}{\sqrt{2}} (-1 + 1) \right|^2 = \boxed{0}$$

no medimos $-\hbar$

$$P_{|V_3, 0\rangle} = \left| \frac{1}{\sqrt{2}} \frac{1}{\sqrt{2}} (-\langle 1,1| + \langle 1,-1|) (|1,1\rangle - |1,-1\rangle) \right|^2 = \left| \frac{1}{2} (-1 - 1) \right|^2 = \boxed{1}$$

Medimos Cero

c) Dada $|\psi\rangle = \frac{1}{\sqrt{2}} |1,1\rangle - \frac{1}{\sqrt{2}} |1,-1\rangle$

medimos $\hat{L}_y \Rightarrow$ mirando la matriz el Av. Corresp es el $|1,1\rangle \Rightarrow$

estamos en $|\psi\rangle$ Después de medir $= |1,1\rangle$

Medimos ahora \hat{L}_y

$$P_{|V_0, 0\rangle} = \left| \frac{1}{\sqrt{2}} (\langle 1,1| + \langle 1,-1|) (|1,1\rangle) \right|^2 = \left| \frac{1}{\sqrt{2}} \right|^2 = \boxed{\frac{1}{2}}$$

$$P_{|V_2, -\hbar\rangle} = \left| \frac{1}{2} (-\langle 1,-1| - \sqrt{2}\langle 1,0| + \langle 1,1|) (|1,1\rangle) \right|^2 = \left| \frac{1}{2} \right|^2 = \boxed{\frac{1}{4}}$$

$$P_{|V_3, +\hbar\rangle} = \left| \frac{1}{2} (-\langle 1,-1| + \sqrt{2}\langle 1,0| + \langle 1,1|) (|1,1\rangle) \right|^2 = \left| \frac{1}{2} \right|^2 = \boxed{\frac{1}{4}}$$

| | | | |
|-----------------|----------|-----------|---------------|
| Podemos obtener | 0 | con prob. | $\frac{1}{2}$ |
| | $-\hbar$ | con prob. | $\frac{1}{4}$ |
| | $+\hbar$ | con prob. | $\frac{1}{4}$ |

15.

$|j, j\rangle$ hasta términos de orden ϵ^2

$$\mathcal{D}_{i\epsilon} |j, j\rangle = e^{-\frac{i\vec{J}_y \epsilon}{\hbar}} |j, j\rangle = e^{-\frac{iJ_y \epsilon}{\hbar}} |j, j\rangle$$

$$\mathcal{D}_{i\epsilon} \approx 1 + \frac{-iJ_y \epsilon}{\hbar} + \frac{(-iJ_y \epsilon)^2}{2\hbar^2} \Rightarrow$$

$$|j, j\rangle_R \approx \left(1 - \frac{iJ_y \epsilon}{\hbar} + \frac{J_y^2 \epsilon^2}{2\hbar^2} \right) |j, j\rangle$$

$$|j, j\rangle_R \approx |j, j\rangle - \frac{i\epsilon}{\hbar} J_y |j, j\rangle + \frac{\epsilon^2}{2\hbar^2} J_y^2 |j, j\rangle$$

$$|j, j\rangle_R \approx |j, j\rangle - \frac{i\epsilon}{\hbar} \left(\frac{J_+ - J_-}{2i} \right) |j, j\rangle + \frac{\epsilon^2}{2\hbar^2} \left(\frac{J_+ - J_-}{4i^2} \right) |j, j\rangle$$

$$\approx |j, j\rangle - \frac{\epsilon}{2\hbar} (J_+ - J_-) |j, j\rangle - \frac{\epsilon^2}{8\hbar^2} (J_+ J_+ - J_+ J_- - J_- J_+ + J_- J_-) |j, j\rangle$$

$$\approx |j, j\rangle - \frac{\epsilon}{2\hbar} (-\hbar \sqrt{j(j+1) - j(j-1)}) |j, j-1\rangle - \frac{\epsilon^2}{8\hbar^2} \left([J_+ - J_-] J_+ |j, j\rangle \right)$$

$$+ (-J_+ + J_-) J_- |j, j\rangle$$

$$\downarrow (-J_+ + J_-) (\hbar \sqrt{j(j+1) - j(j-1)}) |j, j-1\rangle =$$

$$= + \sqrt{j(j+1) - j(j-1)} \hbar^2 \sqrt{j(j+1) - (j-1)[(j-1) - 1]} |j, j-2\rangle - \sqrt{j(j+1) - j(j-2)} \hbar^2 \sqrt{j(j+1) - (j-1)[(j-1) + 1]} |j, j\rangle$$

$$\approx |j, j\rangle + \frac{\epsilon \hbar}{2\hbar} \sqrt{j^2} |j, j-1\rangle - \frac{\epsilon^2}{8\hbar^2} \hbar^2 (\sqrt{j^2} \sqrt{2j-1} |j, j-2\rangle + \sqrt{j^2} \sqrt{2j} |j, j\rangle)$$

$$\approx |j, j\rangle + \frac{\epsilon \sqrt{j}}{\sqrt{2}} |j, j-1\rangle - \frac{\epsilon^2 \sqrt{2j-1}}{4} |j, j-2\rangle - \frac{\epsilon^2 j}{4} |j, j\rangle$$

$$|j, j\rangle_R \underset{\text{orden } \epsilon^2}{=} \left(1 - \frac{\epsilon^2 j}{4} \right) |j, j\rangle + \frac{\epsilon \sqrt{j}}{\sqrt{2}} |j, j-1\rangle - \frac{\epsilon^2 \sqrt{2j-1}}{4} |j, j-2\rangle$$

$$|\langle j, j | j, j\rangle_R|^2 = |\langle j, j | \left(1 - \frac{\epsilon^2 j}{4} \right) |j, j\rangle|^2 = \left| 1 - \frac{\epsilon^2 j}{4} \right|^2 = 1 - \frac{2\epsilon^2 j}{4} + \frac{\epsilon^4 j^2}{16}$$

solo sobreviven los términos en $|j, j\rangle$

$$P_{|j, j\rangle} \approx 1 - \frac{\epsilon^2 j}{2} + O(\epsilon^4)$$

16.

$$(G_i)_{jk} = -i\hbar \epsilon_{ijk}$$

$$G_1 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i\hbar \\ 0 & i\hbar & 0 \end{pmatrix}$$

$$G_2 = \begin{pmatrix} 0 & 0 & i\hbar \\ 0 & 0 & 0 \\ i\hbar & 0 & 0 \end{pmatrix}$$

$$G_3 = \begin{pmatrix} 0 & -i\hbar & 0 \\ i\hbar & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$[G_i, G_j] = i\hbar G_k \epsilon_{ijk}$$

$$G_1 G_2 - G_2 G_1 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i\hbar \\ 0 & i\hbar & 0 \end{pmatrix} \begin{pmatrix} 0 & 0 & i\hbar \\ 0 & 0 & 0 \\ -i\hbar & 0 & 0 \end{pmatrix} - \begin{pmatrix} 0 & 0 & i\hbar \\ 0 & 0 & 0 \\ -i\hbar & 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i\hbar \\ 0 & i\hbar & 0 \end{pmatrix}$$

$$= \begin{pmatrix} 0 & 0 & 0 \\ -i\hbar^2 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} - \begin{pmatrix} 0 & (i\hbar)^2 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} = \begin{pmatrix} 0 & -i\hbar^2 & 0 \\ -i\hbar^2 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$[G_1, G_2] = i\hbar G_3 = i\hbar \begin{pmatrix} 0 & -i\hbar & 0 \\ i\hbar & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} = \begin{pmatrix} 0 & \hbar^2 & 0 \\ -\hbar^2 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \text{satisfies}$$

$$[G_i, G_m]_{jk} = (-i\hbar \epsilon_{ijk})(-i\hbar \epsilon_{mkj}) - (-i\hbar \epsilon_{mjk})(-i\hbar \epsilon_{kij})$$

$$= (-i\hbar) [\epsilon_{ijk} \epsilon_{mkj} - \epsilon_{mjk} \epsilon_{kij}]$$

$$[\epsilon_{jkl} \epsilon_{jmk} - \epsilon_{jkm} \epsilon_{jlk}]$$

$$[\delta_{km} \delta_{ek} - \delta_{kk} \delta_{em} - (\delta_{ke} \delta_{mk} - \delta_{kk} \delta_{me})]$$

$$[\delta_{km} \delta_{ek} - \delta_{ke} \delta_{mk}] \quad k=l=m$$

17.

$V(r) \rightarrow$ H. esféricamente simétrica

$$\Psi(x) = (x+y+3z) f(r)$$

$$\Psi(r, \theta, \phi) = r \cdot (\cos \phi \cdot \sin \theta + \sin \phi \cdot \sin \theta + 3 \cdot \cos \theta) f(r)$$

$$\Psi(r, \theta, \phi) = \underbrace{r \cdot f(r)}_{= R(r)} \cdot (\cos \phi \cdot \sin \theta + \sin \phi \cdot \sin \theta + 3 \cdot \cos \theta) = \langle r, \theta, \phi | n, l, m \rangle$$

$$\Psi(r, \theta, \phi) = R(r) \cdot \sum_l \sum_{m=-l}^l Y_l^m(\theta, \phi)$$

$\sum_l \sum_m \langle r, \theta, \phi | n, l, m \rangle \langle n, l, m | \Psi \rangle = Y_l^m(\theta, \phi)$

queremos ver hasta donde va l

a)

$$\frac{\sqrt{3}}{\sqrt{8\pi}} \cdot \sin \theta \cdot (\cos \phi + i \cdot \sin \phi) = Y_1^1$$

$$\frac{\sqrt{3}}{\sqrt{8\pi}} \cdot \sin \theta \cdot (\cos \phi - i \cdot \sin \phi) = Y_1^{-1}$$

$$\rightarrow Y_1^1 - Y_1^{-1} = 2 \frac{\sqrt{3}}{\sqrt{8\pi}} \sin \theta \cdot \cos \phi$$

$$Y_1^1 + Y_1^{-1} = -2i \frac{\sqrt{3}}{\sqrt{8\pi}} \sin \theta \cdot \sin \phi$$

$$\frac{\sqrt{3}}{\sqrt{4\pi}} \cdot \cos \theta = Y_1^0$$

$$3 \cdot \cos \theta = \sqrt{3 \cdot 4\pi} Y_1^0 = \sqrt{12\pi} Y_1^0$$

$$\Psi(r, \theta, \phi) = R(r) \cdot \left[\sqrt{\frac{8\pi}{3}} \frac{(Y_1^1 - Y_1^{-1})}{2} - \sqrt{\frac{8\pi}{3}} \frac{1}{2i} (Y_1^1 + Y_1^{-1}) + \sqrt{12\pi} Y_1^0 \right]$$

$$\left[\sqrt{\frac{2\pi}{3}} Y_1^1 - \sqrt{\frac{2\pi}{3}} Y_1^{-1} + i \sqrt{\frac{2\pi}{3}} Y_1^1 + i \sqrt{\frac{2\pi}{3}} Y_1^{-1} + \sqrt{12\pi} Y_1^0 \right]$$

$$\left[\sqrt{\frac{2\pi}{3}} (1+i) Y_1^1 + \sqrt{\frac{2\pi}{3}} (-1+i) Y_1^{-1} + \sqrt{12\pi} Y_1^0 \right]$$

$$L^2 \left(\quad \right) + L^2 \left(\quad \right) + L^2 \left(\sqrt{12\pi} Y_1^0 \right)$$

$$\left[2\hbar^2 \left(\quad \right) Y_1^1 + 2\hbar^2 \left(\quad \right) Y_1^{-1} + 2\hbar^2 \left(\quad \right) Y_1^0 \right]$$

\Rightarrow es autofunción de L^2 con autovalor $\boxed{2\hbar^2}$

con $\boxed{l=1}$

b)

$$m=0 \quad P_{m=0} = |\Psi_{m=0}(r, \theta, \phi)|^2$$

$$P_{m=1} = |\Psi_{m=1}(r, \theta, \phi)|^2$$

$$P_{m=-1} = |\Psi_{m=-1}(r, \theta, \phi)|^2$$

Como la parte radial es la misma solo necesitamos normalizar con los $Y_l^m(\theta, \phi) \rightarrow$

$$1 = |\langle Y_1^0 | Y_1^0 \rangle|^2 + |\langle Y_1^1 | Y_1^1 \rangle|^2 + |\langle Y_1^{-1} | Y_1^{-1} \rangle|^2$$

$$\left| \sqrt{\frac{2\pi}{3}} (1+i) \right|^2 + \left| \sqrt{\frac{2\pi}{3}} (-1+i) \right|^2 + \left| \sqrt{12\pi} \right|^2 = 1$$

$$\alpha^2 \left(\left| \sqrt{\frac{2\pi}{3}} \sqrt{2} \right|^2 + \left| \sqrt{\frac{2\pi}{3}} \right|^2 + 12\pi \right) = 1$$

$$\alpha^2 \left(\frac{2\pi}{3} \cdot 2 + \frac{2\pi}{3} + 12\pi \right) = 1$$

$$\alpha^2 \left(\frac{44}{3} \pi \right) = 1$$

$$\alpha = \sqrt{\frac{3}{44\pi}} \rightarrow$$

$$P_{(m=0)} = \frac{3}{44\pi} \cdot \frac{3}{11} = \boxed{\frac{9}{11}}$$

$$P_{(m=1)} = \frac{1}{3} \cdot \frac{3}{11} = \boxed{\frac{1}{11}}$$

$$P_{(m=-1)} = \boxed{\frac{1}{11}}$$

c)

$$H \Psi(x) = E \Psi(x)$$

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + V(r) \right] \Psi(x) - E \Psi(x) = 0$$

$$-\frac{\hbar^2}{2m} \nabla_r^2 \Psi(x) - \frac{\hbar^2}{2m} \nabla_{\theta, \phi}^2 \Psi(x) + V(r) \Psi(x) - E \Psi(x) = 0$$

$$-\frac{\hbar^2}{2m} \nabla_r^2 R(r) \cdot f(\theta, \phi) - \frac{\hbar^2}{2m} \nabla_{\theta, \phi}^2 R(r) \cdot f(\theta, \phi) + V(r) \cdot R(r) \cdot f(\theta, \phi) - E R(r) \cdot f(\theta, \phi) = 0$$

$$+ \frac{L^2}{2m r^2} R(r) \cdot f(\theta, \phi)$$

$$+ \frac{l(l+1)\hbar^2}{2m r^2} R(r) \cdot f(\theta, \phi)$$

$$\boxed{-\frac{\hbar^2}{2m} \nabla_r^2 R(r) + \left[V(r) - E + \frac{l(l+1)\hbar^2}{2m r^2} \right] R(r) = 0}$$

Resolviendo esta ecuación puede hallarse $V(r)$

19.

a) $j=1$

$$\langle 1, m' | J_y | 1, m \rangle$$

↓

$$\langle 1, m' | \frac{J_+ - J_-}{2i} | 1, m \rangle$$

$$\langle 1, 0 | \frac{J_+ - J_-}{2i} | 1, 0 \rangle = \langle 1, 0 | (\# | 1, 1 \rangle - \# | 1, -1 \rangle) = 0$$

$$\langle 1, 0 | \frac{J_+ - J_-}{2i} | 1, 1 \rangle = \langle 1, 0 | (\# | 1, 2 \rangle - \hbar\sqrt{2} | 1, 0 \rangle) = -\frac{\hbar\sqrt{2}}{2i}$$

$$\langle 1, 0 | \frac{J_+ - J_-}{2i} | 1, -1 \rangle = \langle 1, 0 | (\hbar\sqrt{2} | 1, 0 \rangle - \# | 0 \rangle) = \frac{\hbar\sqrt{2}}{2i}$$

$$\langle 1, 1 | \frac{J_+ - J_-}{2i} | 1, 0 \rangle = \langle 1, 1 | (\hbar\sqrt{2} | 1, 1 \rangle - \# | 1, -1 \rangle) = \frac{\hbar\sqrt{2}}{2i}$$

$$\langle 1, 1 | \frac{J_+ - J_-}{2i} | 1, 1 \rangle = \langle 1, 1 | (\# | 0 \rangle - \# | 1, 0 \rangle) = 0$$

$$\langle 1, 1 | \frac{J_+ - J_-}{2i} | 1, -1 \rangle = \langle 1, 1 | (\# | 1, 0 \rangle - \# | 0 \rangle) = 0$$

$$\langle 1, -1 | \frac{J_+ - J_-}{2i} | 1, 0 \rangle = \langle 1, -1 | (\# | 1, 1 \rangle - \hbar\sqrt{2} | 1, -1 \rangle) = \frac{\hbar\sqrt{2}}{2i}$$

$$\langle 1, -1 | \frac{J_+ - J_-}{2i} | 1, 1 \rangle = \langle 1, -1 | (\# | 0 \rangle - \# | 1, 0 \rangle) = 0$$

$$\langle 1, -1 | \frac{J_+ - J_-}{2i} | 1, -1 \rangle = \langle 1, -1 | (\# | 1, 0 \rangle - \# | 0 \rangle) = 0$$

$$J_- | 1, 1 \rangle = \hbar\sqrt{2} | 1, 0 \rangle$$

$$J_+ | 1, -1 \rangle = \hbar\sqrt{2} | 1, 0 \rangle$$

$$J_+ | 1, 0 \rangle = \hbar\sqrt{2} | 1, 1 \rangle$$

$$J_- | 1, 0 \rangle = \hbar\sqrt{2} | 1, -1 \rangle$$

$$J_y = \begin{pmatrix} 0 & \frac{\hbar\sqrt{2}}{2i} & 0 \\ \frac{\hbar\sqrt{2}}{2i} & 0 & -\frac{\hbar\sqrt{2}}{2i} \\ 0 & -\frac{\hbar\sqrt{2}}{2i} & 0 \end{pmatrix}$$

$$J_y = \frac{\hbar}{2} \begin{pmatrix} 0 & -\sqrt{2}i & 0 \\ -\sqrt{2}i & 0 & \sqrt{2}i \\ 0 & \sqrt{2}i & 0 \end{pmatrix}$$

b) caso $j=1$

$$e^{-i J_y \beta / \hbar} = \sum_{n=0}^{\infty} \frac{(-i J_y \beta)^n}{n!} \cdot \frac{1}{n!}$$

$$e^{-i J_y \beta / \hbar} \approx 1 + \frac{-i J_y \beta}{\hbar} + \frac{1}{2} \frac{(-i)^2 J_y^2 \beta^2}{\hbar^2} + \frac{1}{3!} \frac{(-i)^3 J_y^3 \beta^3}{\hbar^3} + \frac{1}{4!} \frac{(-i)^4 J_y^4 \beta^4}{\hbar^4}$$

$$= \sum_{n=0}^{\infty} \frac{1}{(2n)!} \frac{(-i)^{2n} J_y^{2n} \beta^{2n}}{\hbar^{2n}} + \sum_{n=0}^{\infty} \frac{(-i)^{2n+1} J_y^{2n+1} \beta^{2n+1}}{\hbar^{2n+1}} \frac{1}{(2n+1)!}$$

$$= \sum_{n=0}^{\infty} \frac{(-1)^n J_y^{2n} (\beta)^{2n}}{(2n)! (\hbar)^{2n}} + \sum_{n=0}^{\infty} \frac{(-1)^n (-i) (J_y \beta)^{2n+1}}{(2n+1)! (\hbar)^{2n+1}}$$

$$(-i)^{2n} = (-1)^n$$

$$\frac{\partial}{\partial i} (-i)^n = (-1)^n$$

$$\begin{aligned}
 J_y &= \left(\frac{\hbar i}{\sqrt{2}}\right) \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \\
 J_y^2 &= \left(\frac{\hbar i}{\sqrt{2}}\right)^2 \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} = \begin{pmatrix} 1 & 0 & -1 \\ 0 & 2 & 0 \\ -1 & 0 & 1 \end{pmatrix} \left(\frac{\hbar i}{\sqrt{2}}\right)^2 = \left(\frac{\hbar^2}{2}\right) \begin{pmatrix} 1 & 0 & -1 \\ 0 & 2 & 0 \\ -1 & 0 & 1 \end{pmatrix} \\
 J_y^3 &= \left(\frac{\hbar i}{\sqrt{2}}\right)^3 \begin{pmatrix} 1 & 0 & -1 \\ 0 & 2 & 0 \\ -1 & 0 & 1 \end{pmatrix} \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} = \begin{pmatrix} 0 & -2 & 0 \\ 2 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \left(\frac{\hbar i}{\sqrt{2}}\right)^3 = 2 \left(\frac{\hbar i}{\sqrt{2}}\right)^2 J_y = 2 J_y \frac{\hbar^2}{2} \\
 J_y^4 &= \left(\frac{\hbar i}{\sqrt{2}}\right)^4 \begin{pmatrix} 1 & 0 & -1 \\ 0 & 2 & 0 \\ -1 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & -1 \\ 0 & 2 & 0 \\ -1 & 0 & 1 \end{pmatrix} = \begin{pmatrix} 2 & 0 & -2 \\ 0 & 4 & 0 \\ -2 & 0 & 2 \end{pmatrix} \left(\frac{\hbar i}{\sqrt{2}}\right)^4 = 2 J_y^2 \left(\frac{\hbar i}{\sqrt{2}}\right)^2 = J_y^2 \hbar^2 \\
 J_y^6 &= \left(\frac{\hbar i}{\sqrt{2}}\right)^6 \begin{pmatrix} 2 & 0 & -2 \\ 0 & 4 & 0 \\ -2 & 0 & 2 \end{pmatrix} \begin{pmatrix} 1 & 0 & -1 \\ 0 & 2 & 0 \\ -1 & 0 & 1 \end{pmatrix} = \begin{pmatrix} 4 & 0 & -4 \\ 0 & 8 & 0 \\ -4 & 0 & 4 \end{pmatrix} \left(\frac{\hbar i}{\sqrt{2}}\right)^6 = 4 J_y^2 \left(\frac{\hbar i}{\sqrt{2}}\right)^4 = J_y^2 \hbar^4 \\
 J_y^5 &= \left(\frac{\hbar i}{\sqrt{2}}\right)^5 \begin{pmatrix} 1 & 0 & -1 \\ 0 & 2 & 0 \\ -1 & 0 & 1 \end{pmatrix} \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} = J_y^3 \cdot 2 \cdot \left(\frac{\hbar i}{\sqrt{2}}\right)^2 = 2 J_y^3 \left(\frac{\hbar i}{\sqrt{2}}\right)^2 = J_y^3 \hbar^2
 \end{aligned}$$

$$= \sum_{n=0}^{\infty} \frac{(-1)^n J_y^{2n}}{(2n)!} \left(\frac{\beta}{\hbar}\right)^{2n} + i \sum_{n=0}^{\infty} \frac{(-1)^n J_y^{2n+1}}{(2n+1)!} \left(\frac{\beta}{\hbar}\right)^{2n+1}$$

$n \geq 1$

$$\begin{aligned}
 J_y^{2n} &= 2^{n-1} J_y \left(\frac{\hbar i}{\sqrt{2}}\right)^{2n-2} \\
 \left(\frac{J_y}{\hbar}\right)^{2n} &= \cancel{2^{n-1}} \frac{J_y^{2n}}{\hbar^{2n}} = \frac{(i^2)^{n-1}}{\cancel{2^{n-1}}} \\
 \left(\frac{J_y}{\hbar}\right)^{2n} &= \frac{J_y^{2n}}{\hbar^{2n}} (-1)^{n-1}
 \end{aligned}$$

$$\begin{aligned}
 J_y^{2n+1} &= 2^n \left(\frac{\hbar i}{\sqrt{2}}\right)^{2n} J_y \\
 \left(\frac{J_y}{\hbar}\right)^{2n+1} &= \cancel{2^n} \frac{J_y^{2n+1}}{\hbar^{2n+1}} = \frac{J_y^{2n+1}}{\hbar^{2n+1}} \\
 \left(\frac{J_y}{\hbar}\right)^{2n+1} &= (-1)^n \left(\frac{J_y}{\hbar}\right)^{2n+1}
 \end{aligned}$$

$n \geq 1$

$$J_y^3 = \left(\frac{\hbar i}{\sqrt{2}}\right)^3 \begin{pmatrix} 1 & 0 & -1 \\ 0 & 2 & 0 \\ -1 & 0 & 1 \end{pmatrix} \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} = \left(\frac{\hbar i}{\sqrt{2}}\right)^4 \cdot 4 J_y^3 = 24 J_y \left(\frac{\hbar i}{\sqrt{2}}\right)^6 = 8 J_y \left(\frac{\hbar i}{\sqrt{2}}\right)^6 = J_y \hbar^4$$

$$= \sum_{n=0}^{\infty} \frac{(-1)^n J_y^{2n}}{(2n)!} \left(\frac{\beta}{\hbar}\right)^{2n} + \sum_{n=0}^{\infty} \frac{(-1)^n J_y^{2n+1}}{(2n+1)!} \left(\frac{\beta}{\hbar}\right)^{2n+1}$$

$$1 + \sum_{n=1}^{\infty} \frac{(-1)^n J_y^{2n}}{(2n)!} \left(\frac{\beta}{\hbar}\right)^{2n}$$

$$1 + \left(\frac{J_y^2}{\hbar^2}\right) \left[\sum_{n=1}^{\infty} \frac{(-1)^n \beta^{2n}}{(2n)!} \right] + i \left(\frac{J_y}{\hbar}\right) \sum_{n=0}^{\infty} \frac{(-1)^n \beta^{2n+1}}{(2n+1)!}$$

$$\begin{aligned}
 (i^2)^n &= (-1)^n \\
 (-1)^n &= (-1)^n \\
 (-1)^n &= (-1)^n
 \end{aligned}$$

c)

$$j=1 \quad e^{-\frac{iJ_y \beta}{\hbar}} = 1 - i \left(\frac{J_y}{\hbar} \right) \sin \beta - \left(\frac{J_y}{\hbar} \right)^2 (1 - \cos \beta)$$

$$d^{j=1}(\beta) = \begin{pmatrix} \langle 1, -1 | e^{-\frac{iJ_y \beta}{\hbar}} | 1, -1 \rangle & \langle 1, -1 | e^{-\frac{iJ_y \beta}{\hbar}} | 1, 0 \rangle & \langle 1, -1 | e^{-\frac{iJ_y \beta}{\hbar}} | 1, 1 \rangle \\ \langle 1, 0 | e^{-\frac{iJ_y \beta}{\hbar}} | 1, -1 \rangle & \langle 1, 0 | e^{-\frac{iJ_y \beta}{\hbar}} | 1, 0 \rangle & \langle 1, 0 | e^{-\frac{iJ_y \beta}{\hbar}} | 1, 1 \rangle \\ \langle 1, 1 | e^{-\frac{iJ_y \beta}{\hbar}} | 1, -1 \rangle & \langle 1, 1 | e^{-\frac{iJ_y \beta}{\hbar}} | 1, 0 \rangle & \langle 1, 1 | e^{-\frac{iJ_y \beta}{\hbar}} | 1, 1 \rangle \end{pmatrix}$$

$$\frac{J_y^2}{\hbar^2} = \begin{pmatrix} -1 & 0 & 1 \\ 1/2 & 0 & -1/2 \\ 0 & 1 & 0 \\ -1/2 & 0 & 1/2 \end{pmatrix} \begin{matrix} \langle 1 \\ \langle 0 \\ \langle 1 \end{matrix}$$

$$\frac{J_y}{\hbar} = \begin{pmatrix} -1 & 0 & 1 \\ 0 & -i/\sqrt{2} & 0 \\ -i/\sqrt{2} & 0 & i/\sqrt{2} \\ 0 & 1/\sqrt{2} & 0 \end{pmatrix} \begin{matrix} \langle 1 \\ \langle 0 \\ \langle 1 \end{matrix}$$

$$d_{11}^{j=1}(\beta) = 1 - \frac{1}{2}(1 - \cos \beta) = 1 - \frac{1}{2} + \frac{\cos \beta}{2}$$

$$d_{12}^{j=1}(\beta) = -i \left(\frac{1}{\sqrt{2}} \right) \sin \beta = -\frac{1}{\sqrt{2}} \sin \beta$$

$$d_{13}^{j=1}(\beta) = -\left(\frac{1}{2} \right) (1 - \cos \beta) =$$

$$d_{21}^{j=1}(\beta) = -i \left(\frac{-i}{\sqrt{2}} \right) \sin \beta = -\frac{1}{\sqrt{2}} \sin \beta$$

$$d_{22}^{j=1}(\beta) = 1 - (1 - \cos \beta) = \cos \beta$$

$$d_{23}^{j=1}(\beta) = -i \left(\frac{i}{\sqrt{2}} \right) \sin \beta = \frac{1}{\sqrt{2}} \sin \beta$$

$$d_{31}^{j=1}(\beta) = -\left(\frac{1}{2} \right) (1 - \cos \beta) =$$

$$d_{32}^{j=1}(\beta) = -i \left(\frac{i}{\sqrt{2}} \right) \sin \beta = \frac{1}{\sqrt{2}} \sin \beta$$

$$d_{33}^{j=1}(\beta) = 1 - \frac{1}{2}(1 - \cos \beta) = 1 - \frac{1}{2} + \frac{\cos \beta}{2}$$

$$d^{j=1}(\beta) = \begin{pmatrix} \frac{1}{2}(1 + \cos \beta) & -\frac{1}{\sqrt{2}} \sin \beta & \frac{1}{2}(1 - \cos \beta) \\ -\frac{1}{\sqrt{2}} \sin \beta & \cos \beta & \frac{1}{\sqrt{2}} \sin \beta \\ \frac{1}{2}(1 - \cos \beta) & \frac{1}{\sqrt{2}} \sin \beta & \left(\frac{1}{2} \right) (1 + \cos \beta) \end{pmatrix}$$