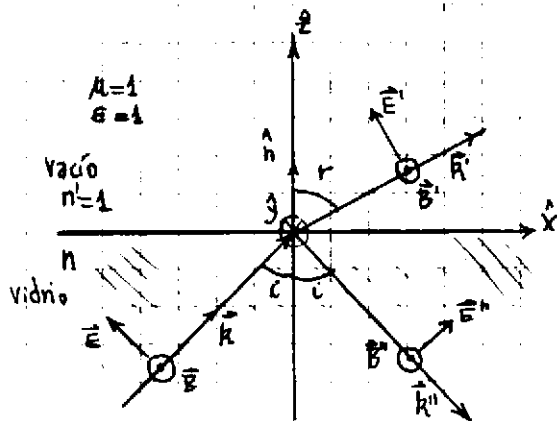


GUÍA 6

I. Problemas

1.



onda l.p. con TM, $n > n' = 1$

$$i > i_0 = \text{asen} \left(\frac{n'}{n} \right) \quad \frac{n'}{n} < 1$$

ángulo límite

$$\text{sen } i > \frac{n'}{n}$$

$$n \cdot \text{sen } i = n' \cdot \text{sen } r \quad (\text{Snell law})$$

$$\text{sen } i > \frac{\text{sen } i}{\text{sen } r} \Rightarrow \text{sen } r > 1$$

$$\therefore r \in \mathbb{C} \Rightarrow$$

$$\text{sen } r > 1 \rightarrow \cos r = \sqrt{-\text{sen}^2 r + 1} = i \sqrt{\text{sen}^2 r - 1}$$

$$\cos r = i \cdot \sqrt{\text{sen}^2 i \left(\frac{n^2}{n'^2} \right) - 1}$$

$$\cos r = i \cdot \sqrt{\left(\frac{\text{sen } i}{\text{sen } i_0} \right)^2 - 1}$$

$$\therefore \cos r \in \mathbb{Im}$$

a.

$$\langle \vec{S} \rangle = \frac{c}{8\pi} \text{Re} \{ \vec{E} \times \vec{H}^* \}$$

$$\langle \vec{S} \rangle = \frac{c}{8\pi} \text{Re} \{ \vec{E} \times \vec{H}^* \}$$

$$\langle \vec{S} \cdot \hat{n} \rangle = \frac{c}{8\pi} \text{Re} \left\{ \frac{c}{\mu' \omega} (\vec{E} \times \vec{k}'^* \cdot \vec{E}^*) \cdot \hat{n} \right\}$$

$$= \frac{c^2}{8\pi \mu' \omega} \text{Re} \left\{ (\vec{k}'^* \cdot \vec{E}^* - \vec{E}^* \cdot (\vec{E} \cdot \vec{k}'^*)) \cdot \hat{n} \right\}$$

$$\langle \vec{S} \cdot \hat{n} \rangle = \frac{c^2}{8\pi \mu' \omega} \text{Re} \{ \vec{k}'^* \cdot |\vec{E}^*|^2 \cdot \hat{n} \}$$

$$\frac{c^2 \cdot |\vec{E}^*|^2}{8\pi \mu' \omega} \text{Re} \{ \vec{k}'^* \cdot \hat{n} \}$$

$$\text{Re} \{ \cos r \cdot k' \} = 0$$

$$\Rightarrow \langle \vec{S} \cdot \hat{n} \rangle = 0$$

Como el ángulo de i es $> i_0 \rightarrow$

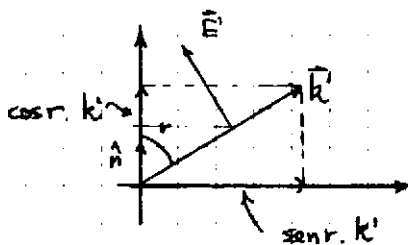
$$\vec{k}' \in \mathbb{C}$$

$$\vec{B}' = \sqrt{\mu' \epsilon'} \frac{\vec{k}' \times \vec{E}'}{k'}$$

Como el modo primario es vacío \therefore

$$\vec{H}' = \frac{\sqrt{\mu' \epsilon'}}{\mu'} \cdot \frac{(\vec{k}' \times \vec{E}')}{k'}$$

$$\vec{H}'^* = \frac{\sqrt{\mu' \epsilon'}}{\mu'} \cdot \frac{(\vec{k}' \times \vec{E}')^*}{\sqrt{\mu' \epsilon'} \cdot \frac{\omega}{c}}$$



No hay flujo del vector de Poynting en la dirección normal a la superficie

b. El ángulo de Brewster es tal que:

$$\text{Brewster } i_B = \arctan\left(\frac{n'}{n}\right)$$

$$(> \text{ límite}) \rightarrow i > \arcsen\left(\frac{n'}{n}\right)$$

Para incidir además con ángulo de Brewster necesitamos i' :

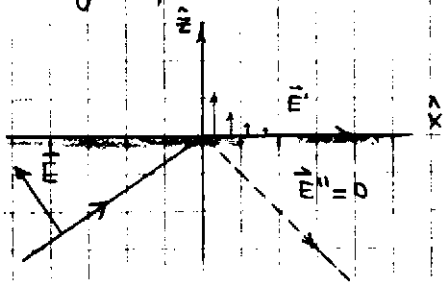
$$\tan i' = \frac{n'}{n} \quad \text{pero} \quad \text{sen } i' > \frac{n'}{n}$$

$$\text{luego quiero } i': \quad \text{sen } i' > \frac{\text{sen } i'}{\text{cos } i'} \Rightarrow \text{cos } i' > 1$$

No hay i' real que verifique esto

Parece lógico que no puedan darse ambas situaciones en simultáneo, porque eso significaría que la interfase se engulle a la onda.

Esto significa que el medio está tragándose constantemente energía de la onda, lo cual parecería violar la conservación.



2.

$$i > i_0 \rightarrow \vec{k} \in \mathbb{C}$$

incide luz LP

$$\vec{E} = (\hat{e}_\parallel E_\parallel + \hat{e}_\perp E_\perp) e^{i(\vec{k} \cdot \vec{x} + \omega t)}$$

$$\vec{E}'' = (\hat{e}_\parallel E''_\parallel + \hat{e}_\perp E''_\perp) e^{i(\vec{k}'' \cdot \vec{x} - \omega t)}$$

* D_n continuo en contorno

$$\epsilon (E_\parallel + E''_\parallel) \cdot \text{sen } i - \epsilon' E''_\parallel \text{sen } r = 0$$

$$\epsilon (E_\perp + E''_\perp) \cdot \sqrt{\frac{\epsilon \mu'}{\epsilon \mu}} \text{sen } r - \epsilon' E''_\perp \text{sen } r = 0$$

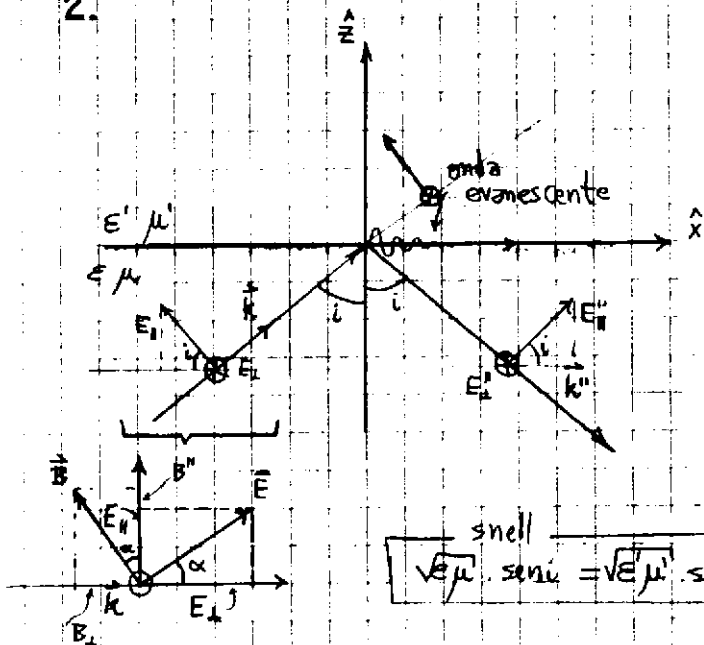
$$\sqrt{\epsilon} (E_\parallel + E''_\parallel) \sqrt{\mu} - \sqrt{\epsilon'} \sqrt{\mu'} E''_\parallel = 0$$

* E_z continuo en contorno

$$E_\perp + E''_\perp - E'_\perp = 0$$

$$E_\perp + E''_\perp = E'_\perp$$

$$\sqrt{\frac{\epsilon}{\mu}} (E \cdot \text{sen } \alpha + E''_\parallel) = \sqrt{\frac{\epsilon'}{\mu'}}$$



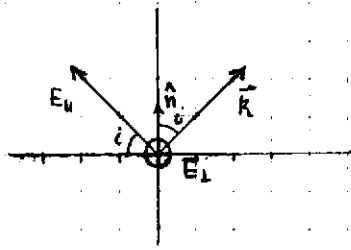
snell
 $\sqrt{\epsilon \mu} \text{sen } i = \sqrt{\epsilon' \mu'} \text{sen } r$

$$\begin{aligned} E_\perp &= E \cdot \text{cos } \alpha \\ E_\parallel &= E \cdot \text{sen } \alpha \\ B_\parallel &= B \cdot \text{cos } \alpha \\ B_\perp &= B \cdot \text{sen } \alpha \end{aligned}$$

* B_n continuo

$$[(\vec{k} \times \vec{E}_0 + \vec{k}'' \times \vec{E}''_0) - (\vec{k}' \times \vec{E}'_0)] \cdot \hat{n} = 0$$

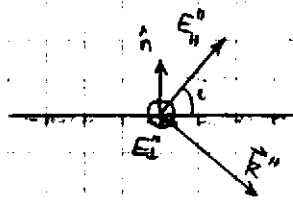
$$[(\vec{k} \times \vec{E}_\perp) \cdot \hat{n} + (\vec{k}'' \times \vec{E}''_\perp) \cdot \hat{n}] - [(\vec{k}' \times \vec{E}'_\perp) \cdot \hat{n}] = 0$$



$$\begin{aligned} \vec{k} \times \vec{E}_0 \cdot \hat{n} &= 0 \\ \vec{k}' \times \vec{E}_1 \cdot \hat{n} &\neq 0 \end{aligned}$$

$$k \cdot E_0 \cos\left(\frac{\pi}{2} - i\right)$$

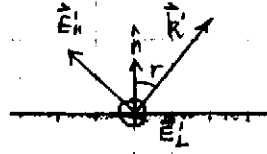
$$k \cdot E_1 \sin i$$



$$\begin{aligned} \vec{k}'' \times \vec{E}_0'' \cdot \hat{n} &= 0 \\ \vec{k}'' \times \vec{E}_1'' \cdot \hat{n} &\neq 0 \end{aligned}$$

$$k'' \cdot E_0'' \cos\left(\frac{\pi}{2} - i\right)$$

$$k'' \cdot E_1'' \sin i$$



$$\begin{aligned} \vec{k}' \times \vec{E}_0' \cdot \hat{n} &= 0 \\ \vec{k}' \times \vec{E}_1' \cdot \hat{n} &\neq 0 \end{aligned}$$

$$k' \cdot E_0' \cos\left(\frac{\pi}{2} - r\right)$$

$$k' \cdot E_1' \sin r$$

$$k \cdot E_0 \sin i + k'' \cdot E_1'' \sin i = k' \cdot E_1' \sin r = k' \cdot E_0' \sin i \sqrt{\frac{\epsilon \mu}{\epsilon' \mu'}}$$

$$\boxed{k'' = k}$$

$$k(E_0 + E_1) = k' \cdot E_0' \sqrt{\frac{\epsilon \mu}{\epsilon' \mu'}}$$

$$\boxed{k' = k \sqrt{\frac{\epsilon' \mu'}{\epsilon \mu}}}$$

* Hz continuo

$$\left[\frac{1}{\mu} (\vec{k} \times \vec{E}_0 + \vec{k}'' \times \vec{E}_0'') - \frac{1}{\mu'} (\vec{k}' \times \vec{E}_0') \right] \times \hat{n} = 0$$

$$\left[-\frac{1}{\mu} \hat{n} \times (\vec{k} \times \vec{E}_0) - \frac{1}{\mu} \hat{n} \times (\vec{k}'' \times \vec{E}_0'') \right] + \left[\hat{n} \times \frac{1}{\mu'} \vec{k}' \times \vec{E}_0' \right] = 0$$

$$-\frac{1}{\mu} \left[\vec{k} (\hat{n} \cdot \vec{E}_0) - \vec{E}_0 (\hat{n} \cdot \vec{k}) \right] - \frac{1}{\mu} \left[\vec{k}'' (\hat{n} \cdot \vec{E}_0'') - \vec{E}_0'' (\hat{n} \cdot \vec{k}'') \right]$$

$$+ \frac{1}{\mu'} \left[\vec{k}' (\hat{n} \cdot \vec{E}_0') - \vec{E}_0' (\hat{n} \cdot \vec{k}') \right] = 0$$

$$-\frac{1}{\mu} \left[\vec{k} \cdot E_0 \cos\left(\frac{\pi}{2} - i\right) - \vec{E}_0 \cdot k \cos i \right] - \frac{1}{\mu} \left[\vec{k}'' \cdot E_0'' \cos\left(\frac{\pi}{2} - i\right) - \vec{E}_0'' \cdot k'' \cos i \right]$$

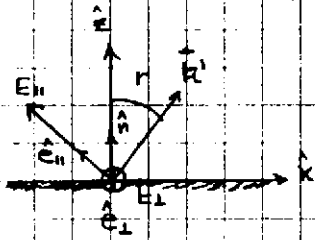
$$+ \frac{1}{\mu'} \left[\vec{k}' \cdot E_0' \cos\left(\frac{\pi}{2} - r\right) - \vec{E}_0' \cdot k' \cos r \right] = 0$$

$$-\frac{\vec{k} \cdot E_0 \sin i}{\mu} + \frac{\vec{E}_0 \cdot k \cos i}{\mu} - \frac{\vec{k}'' \cdot E_0'' \sin i}{\mu} + \frac{\vec{E}_0'' \cdot k'' \cos i}{\mu} + \frac{\vec{k}' \cdot E_0' \sin r}{\mu'} - \frac{\vec{E}_0' \cdot k' \cos r}{\mu'} = 0$$

$$\vec{E}' = (\hat{e}_{||} E'_{||} + \hat{e}_{\perp} E'_{\perp}) \cdot e^{i(\vec{k}' \cdot \vec{x} - \omega t)}$$

$$e^{i[k'(\cos r z + \sin r x) - \omega t]}$$

$$(\hat{e}_{||} E'_{||} + \hat{e}_{\perp} E'_{\perp}) e^{-k' |\cos r| z} \cdot e^{i[k' \sin r x - \omega t]}$$



* Coeficientes de Fresnel

E paralelo

$$\frac{E'_{||}}{E_{||}} = \frac{2 n n' \cos i}{\frac{\mu}{\mu'} n^2 \cos i + n n' \cos r} = \frac{E_{||}}{E \cdot \sin \alpha} = \frac{2 \sqrt{\mu' \mu} E \cdot \cos i}{\frac{\mu}{\mu'} E \cdot \cos i + \sqrt{\mu' \mu} E \cdot \cos r}$$

$$\frac{\sqrt{\mu' \mu} \cdot 2 \sqrt{\mu' \mu} \cdot \cos i}{\sqrt{\mu' \mu} (\sqrt{\mu' \mu} \cdot \cos i + \sqrt{\mu' \mu} \cdot \cos r)}$$

$$E'_{||} = \frac{E \cdot \sin \alpha \cdot 2 \sqrt{\mu' \mu} \cdot \cos i (\sqrt{\mu' \mu} \cos i - \sqrt{\mu' \mu} \cos r)}{\mu' \cos^2 i + \mu \cos^2 r}$$

$$\frac{E'_{||}}{E_{||}} = \frac{\frac{\mu}{\mu'} n^2 \cos i - n n' \cos r}{\frac{\mu}{\mu'} n^2 \cos i + n n' \cos r} = \frac{\left[\frac{\mu}{\mu'} n^2 \cos i - n n' \cos r \right]^2}{\left[\frac{\mu}{\mu'} n^2 \cos i \right]^2 + \left[n n' \cos r \right]^2}$$

$$= \frac{\left[\frac{\mu}{\mu'} n^2 \cos i \right]^2 - 2 \left(\frac{\mu}{\mu'} n^2 \cos i \right) (n n' \cos r) - (n n' \cos r)^2}{\left[\frac{\mu}{\mu'} n^2 \cos i \right]^2 + \left[n n' \cos r \right]^2}$$

$$E'_{||} = \left(\frac{A^2 - B^2}{A^2 + B^2} - i \frac{2AB}{A^2 + B^2} \right) \cdot \sin \alpha \cdot E$$

E perpendicular

$$\frac{E'_{\perp}}{E_{\perp}} = \frac{2 n \cos i}{n \cos i + \frac{\mu}{\mu'} \sqrt{n^2 - n'^2} \sin^2 i} = \frac{E'_{\perp}}{E \cos \alpha} = \frac{2 \sqrt{\mu' \mu} \cdot \cos i}{\sqrt{\mu' \mu} \cdot \cos i + \frac{\mu}{\mu'} \sqrt{\mu' \mu} \cdot \cos r}$$

$$E'_{\perp} = \frac{2 \cos i \cdot \cos \alpha \cdot E}{\cos i + \frac{\mu \sqrt{\mu' \mu}}{\mu' \mu} \cdot \cos r}$$

$$E'_{\perp} = \frac{2 \cos i \cdot \cos \alpha \cdot E (\cos i - |\cos r| \cdot i \cdot \left[\frac{\mu' \mu}{\mu \mu'} \right])}{\cos^2 i + \frac{\mu \mu'}{\mu \mu'} \cdot |\cos r|^2}$$

$$\frac{E'_{\perp}}{E_{\perp}} = \frac{\sqrt{\mu' \mu} \cos i - \frac{\mu}{\mu'} \sqrt{\mu' \mu} \cos r}{\sqrt{\mu' \mu} \cos i + \frac{\mu}{\mu'} \sqrt{\mu' \mu} \cos r} = \frac{\sqrt{\mu' \mu} \cos i - \frac{\mu}{\mu'} \sqrt{\mu' \mu} \cos r}{\sqrt{\mu' \mu} \cos i + \frac{\mu}{\mu'} \sqrt{\mu' \mu} \cos r}$$

$$= \frac{\cos i - \frac{\mu \mu'}{\mu \mu'} \cos r}{\cos i + \frac{\mu \mu'}{\mu \mu'} \cos r}$$

$$\frac{E'_{\perp}}{E_{\perp}} = \frac{\cos^2 i - 2 \cos i \cdot |\cos r| \cdot i \cdot \frac{\mu \mu'}{\mu \mu'} - \frac{\mu \mu'}{\mu \mu'} \cos^2 r}{\cos^2 i + |\cos r|^2 \cdot \frac{\mu \mu'}{\mu \mu'}}$$

$$E_{\perp}'' = \frac{E \cdot \cos \alpha \cdot (\cos^2 i - \frac{\mu \epsilon'}{\epsilon \mu'} |\cos r|^2)}{\cos^2 i + |\cos r|^2 \frac{\mu \epsilon'}{\epsilon \mu'}} - \frac{i E \cdot \cos \alpha \cdot 2 \cos i |\cos r| \sqrt{\frac{\mu \epsilon'}{\mu' \epsilon}}}{\cos^2 i + |\cos r|^2 \frac{\mu \epsilon'}{\epsilon \mu'}}$$

$$|E_{\perp}''| = \sqrt{\frac{[(\cos^2 i - \frac{\mu \epsilon'}{\epsilon \mu'} |\cos r|^2) \cdot E \cos \alpha]^2 + [E \cos \alpha \cdot 2 \cos i |\cos r| \sqrt{\frac{\mu \epsilon'}{\mu' \epsilon}}]^2}{[\cos^2 i + |\cos r|^2 \frac{\mu \epsilon'}{\epsilon \mu'}]^2}}$$

$$= \frac{E \cdot \cos \alpha \cdot \sqrt{(\cos^2 i + \frac{\mu \epsilon'}{\epsilon \mu'} |\cos r|^2) - 2 \cos i \cdot \frac{\mu \epsilon'}{\epsilon \mu'} |\cos r|^2 + 4 \cos^2 i |\cos r|^2 \frac{\mu \epsilon'}{\epsilon \mu'}}}{\cos^2 i + |\cos r|^2 \frac{\mu \epsilon'}{\epsilon \mu'}}$$

$$|E_{\perp}''| = \frac{E \cdot \cos \alpha \cdot (\cos^2 i + |\cos r|^2 \frac{\mu \epsilon'}{\epsilon \mu'})}{\cos^2 i + |\cos r|^2 \frac{\mu \epsilon'}{\epsilon \mu'}} \rightarrow |E_{\perp}''| = E \cdot \cos \alpha$$

$$\varphi_{\perp} = \arctan \left\{ \frac{\text{Im } E_{\perp}''}{\text{Re } E_{\perp}''} \right\} \Rightarrow$$

$$\varphi_{\perp} = \arctan \left\{ \frac{-E \cdot \cos \alpha \cdot 2 \cos i \cdot |\cos r| \cdot \sqrt{\frac{\mu \epsilon'}{\mu' \epsilon}}}{E \cdot \cos \alpha \cdot (\cos^2 i - |\cos r|^2 \frac{\mu \epsilon'}{\epsilon \mu'})} \right\}$$

$$|E_{\parallel}''| = E \cdot \sin \alpha \cdot \sqrt{\frac{\left[\left(\frac{\mu n'^2 \cos i}{\mu'} \right) - [n n' |\cos r|^2] \right]^2 + \left(2 \cdot \frac{\mu n'^2 \cos i \cdot n n' \cdot |\cos r|}{\mu'} \right)^2}{\sqrt{\left[\left(\frac{\mu n'^2 \cos i}{\mu'} \right)^2 + (n n' |\cos r|)^2 \right]^2}}$$

$$|E_{\parallel}''| = E \cdot \sin \alpha$$

$$\varphi_{\parallel} = \arctan \left\{ \frac{\text{Im } E_{\parallel}''}{\text{Re } E_{\parallel}''} \right\} \Rightarrow$$

$$\varphi_{\parallel} = \arctan \left\{ - \frac{2 \cdot \frac{\mu n'^2 \cos i \cdot n n' |\cos r|}{\mu}}{\left[\frac{\mu n'^2 \cos i}{\mu'} \right]^2 - [n n' |\cos r|]^2} \right\}$$

$$\varphi_{\parallel} = \arctan \left\{ \frac{-2 \cdot \frac{\mu n'^2 \cos i}{\mu} \cdot \sqrt{\mu \epsilon'} \cdot \mu' \epsilon' |\cos r| \cdot \cos i}{\left(\frac{\mu n'^2 \cos i}{\mu'} \right)^2 - (\sqrt{\mu \epsilon'} \mu' \epsilon' |\cos r|)^2} \right\}$$

* Ángulo r

$$\cos r = i \cdot |\cos r| =$$

$$\cos^2 r = 1 - \sin^2 r$$

$$\cos r = i \cdot \sqrt{\sin^2 r - 1}$$

$$\cos r = i \cdot \sqrt{\frac{\sin^2 i \cdot \frac{\epsilon \mu}{\epsilon' \mu'} - 1}{\epsilon' \mu'}}$$

$$|\cos r|^2 = \left(\frac{\sin^2 i \cdot \frac{\epsilon \mu}{\epsilon' \mu'} - 1}{\epsilon' \mu'} \right)$$

$$\varphi_{\parallel} = \arctan \left\{ \frac{-2 \cos i \cdot \mu \epsilon' \sqrt{\mu \epsilon'} \cdot \mu' \epsilon' \cdot \sqrt{\sin^2 i \cdot \frac{\epsilon \mu}{\epsilon' \mu'} - 1}}{(\mu \epsilon')^2 \cos^2 i - \mu \epsilon' \mu' \epsilon' \left(\frac{\sin^2 i \cdot \frac{\epsilon \mu}{\epsilon' \mu'} - 1}{\epsilon' \mu'} \right)} \right\}$$

$$\varphi_{\parallel} = \arctan \left\{ \frac{-2 \cos i \cdot \mu \cdot \sqrt{\mu \epsilon'} \cdot \mu' \epsilon' \cdot \sqrt{\sin^2 i \cdot \frac{\epsilon \mu}{\epsilon' \mu'} - 1}}{\mu^2 \epsilon'^2 \cos^2 i - \mu^2 \epsilon'^2 \sin^2 i + \mu \epsilon' \mu' \epsilon'}$$

$$\varphi_{\parallel} = \arctan \left\{ \frac{-2 \cos i \cdot \sqrt{\frac{\epsilon \mu'}{\mu \epsilon'}} \cdot \sqrt{\sin^2 i \cdot \frac{\epsilon \mu}{\epsilon' \mu'} - 1}}{\cos^2 i - \left(\frac{\epsilon}{\epsilon'} \right)^2 \sin^2 i + \frac{\epsilon \mu'}{\mu \epsilon'}}$$

$$\vec{E}^{\text{refl}} = \left(\hat{e}_{\parallel} E \sin \alpha \cdot e^{i\varphi_{\parallel}} + \hat{e}_{\perp} E \cos \alpha \cdot e^{i\varphi_{\perp}} \right) \cdot e^{i(\vec{k} \cdot \vec{x} - \omega t)}$$

$$\varphi_{\perp} = \text{atan} \left\{ \frac{-2 \cdot \cos i \cdot \sqrt{\sin^2 i \cdot (\epsilon \mu / \epsilon' \mu') - 1} \cdot \sqrt{\frac{\epsilon' \mu'}{\epsilon \mu}}}{\cos^2 i - \frac{\epsilon' \mu'}{\epsilon \mu} \cdot \left(\frac{\sin^2 i \cdot \epsilon \mu}{\epsilon' \mu'} - 1 \right)} \right\}$$

$$\varphi_{\perp} = \text{atan} \left\{ \frac{-2 \cos i \cdot \sqrt{\sin^2 i \cdot (\mu / \mu')^2 - \epsilon' \mu / \epsilon \mu'}}{\cos^2 i + \frac{\epsilon' \mu}{\epsilon \mu'} - \frac{\sin^2 i \cdot (\mu / \mu')^2}{\mu'}}$$

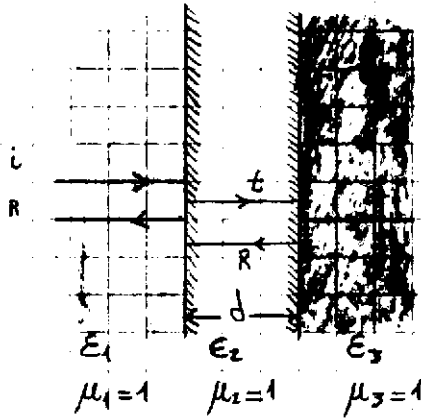
$$\varphi_{\parallel} = \text{atan} \left\{ \frac{-2 \cdot \cos i \cdot \sqrt{\sin^2 i \cdot (\epsilon / \epsilon')^2 - \epsilon' / \epsilon \mu}}{\cos^2 i + \frac{\epsilon \mu'}{\epsilon' \mu} - \sin^2 i \cdot \left(\frac{\epsilon'}{\epsilon} \right)^2} \right\}$$

si suponemos frecuencias ópticas \Rightarrow

$$\varphi_{\perp} = \text{atan} \left\{ \frac{-2 \cdot \cos i \cdot \sqrt{\sin^2 i - \epsilon' / \epsilon}}{\cos^2 i + \frac{\epsilon'}{\epsilon} - \sin^2 i} \right\}$$

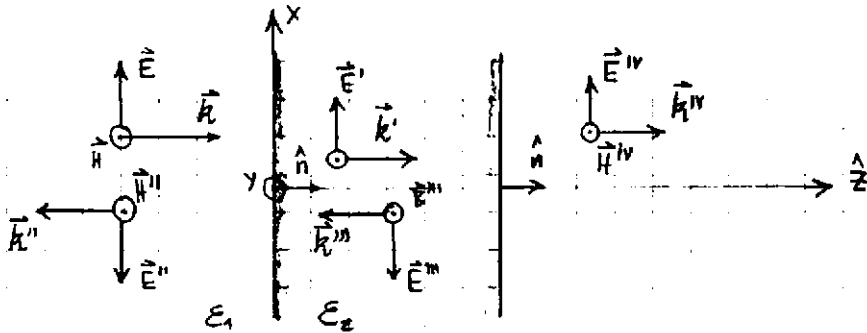
$$\varphi_{\parallel} = \text{atan} \left\{ \frac{-2 \cdot \cos i \cdot \sqrt{\sin^2 i \cdot (\epsilon / \epsilon')^2 - \epsilon' / \epsilon}}{\cos^2 i + \frac{\epsilon}{\epsilon'} - \sin^2 i \cdot \left(\frac{\epsilon'}{\epsilon} \right)^2} \right\}$$

3.



$$\vec{E}(\vec{x}, t) = \vec{E}_0 \cdot e^{i(\vec{k} \cdot \vec{x} - \omega t)}$$

a.



* \vec{E}_t continuo ($z=0$)

$$\begin{aligned} [(\vec{E} + \vec{E}'') - (\vec{E}' + \vec{E}''')] \times \hat{n} &= 0 \\ (-E_y \hat{y} + E''_y \hat{y}) - (-E'_y \hat{y} + E'''_y \hat{y}) &= 0 \\ -E_0 + E''_0 &= -E'_0 + E'''_0 \end{aligned}$$

* \vec{H}_t continuo ($z=0$)

$$[(\vec{k} \times \vec{E}_0 + \vec{k}'' \times \vec{E}''_0) - (\vec{k}' \times \vec{E}'_0 + \vec{k}''' \times \vec{E}'''_0)] \times \hat{n} = 0$$

x)

$$\begin{aligned} k E_0 + k'' E''_0 - k' E'_0 - k''' E'''_0 &= 0 \\ \frac{\sqrt{\epsilon_1 \mu_1} \omega}{c} E_0 + \frac{\sqrt{\epsilon_1 \mu_1} \omega}{c} E''_0 - \frac{\sqrt{\epsilon_2 \mu_2} \omega}{c} E'_0 - \frac{\sqrt{\epsilon_2 \mu_2} \omega}{c} E'''_0 &= 0 \\ \sqrt{\epsilon_1} E_0 + \sqrt{\epsilon_1} E''_0 - \sqrt{\epsilon_2} E'_0 - \sqrt{\epsilon_2} E'''_0 &= 0 \rightarrow \sqrt{\epsilon_1} (E_0 + E''_0) = \sqrt{\epsilon_2} (E'_0 + E'''_0) \end{aligned}$$

* \vec{E}_t continuo ($z=d$)

$$[(\vec{E}'_0 e^{ik'd} + \vec{E}'''_0 e^{-ik'''d}) - \vec{E}''_0 e^{ik''d}] \times \hat{n} = 0$$

-y)

$$\begin{aligned} -E'_0 e^{ik'd} + E'''_0 e^{-ik'''d} - (-E''_0 e^{ik''d}) &= 0 \\ -E'_0 e^{i\sqrt{\epsilon_2 \mu_2} \omega d / c} + E'''_0 e^{-i\sqrt{\epsilon_2 \mu_2} \omega d / c} &= -E''_0 e^{i\sqrt{\epsilon_2 \mu_2} \omega d / c} \\ -E'_0 + E'''_0 e^{-i\sqrt{\epsilon_2} \omega d / c} &= -E''_0 e^{i\sqrt{\epsilon_2} \omega d / c} \\ -E'_0 + E'''_0 e^{-i\sqrt{\epsilon_2} \omega d / c} &= -E''_0 e^{i(\epsilon_3 - \sqrt{\epsilon_2}) \omega d / c} \\ -E'_0 + E'''_0 e^{-i\sqrt{\epsilon_2} \omega d / c} &= -E''_0 e^{(\epsilon_3 - \sqrt{\epsilon_2}) \omega d / c} \end{aligned}$$

* \vec{H}_t continuo ($z=d$)

$$\begin{aligned}
 & \left[(\vec{k}^I \times \vec{E}^I + \vec{k}^{III} \times \vec{E}^{III}) - (\vec{k}^{IV} \times \vec{E}^{IV}) \right] \times \hat{n} = 0 \\
 & k^I E_0^I e^{i k^I d} + k^{III} E_0^{III} e^{-i k^{III} d} - k^{IV} E_0^{IV} e^{i k^{IV} d} = 0 \\
 & \sqrt{\epsilon_2} \mu_2 \frac{\omega}{c} E_0^I e^{i \sqrt{\epsilon_2} \mu_2 \frac{\omega}{c} d} + \sqrt{\epsilon_2} \mu_2 \frac{\omega}{c} E_0^{III} e^{-i \sqrt{\epsilon_2} \mu_2 \frac{\omega}{c} d} = \sqrt{\epsilon_3} \mu_3 \frac{\omega}{c} E_0^{IV} e^{i \sqrt{\epsilon_3} \mu_3 \frac{\omega}{c} d} \\
 & \sqrt{\epsilon_2} \cdot E_0^I \cdot e^{i \sqrt{\epsilon_2} \frac{2 \omega d}{c}} + \sqrt{\epsilon_2} \cdot E_0^{III} \cdot e^{-i \sqrt{\epsilon_2} \frac{2 \omega d}{c}} = \sqrt{\epsilon_3} \cdot E_0^{IV} \cdot e^{i \sqrt{\epsilon_3} \frac{2 \omega d}{c}} \\
 & \sqrt{\epsilon_2} \left(E_0^I + E_0^{III} e^{-i \sqrt{\epsilon_2} \frac{2 \omega d}{c}} \right) = \sqrt{\epsilon_3} e^{i \sqrt{\epsilon_3} \frac{2 \omega d}{c}} E_0^{IV}
 \end{aligned}$$

- 1 $-E_0 + E_0^{IV} = -E_0^I + E_0^{III}$ Necesito $E_0^{IV} = 0$
- 2 $\sqrt{\epsilon_1} (E_0 + E_0^{IV}) = \sqrt{\epsilon_2} (E_0^{III} + E_0^I)$
- 3 $E_0 - E_0^{IV} e^{i \sqrt{\epsilon_2} \frac{2 \omega d}{c}} = E_0^{III} e^{i \sqrt{\epsilon_3} \frac{2 \omega d}{c}} - E_0^I$
- 4 $\sqrt{\epsilon_2} (E_0 + E_0^{IV} e^{-i \sqrt{\epsilon_2} \frac{2 \omega d}{c}}) = \sqrt{\epsilon_3} e^{i \sqrt{\epsilon_3} \frac{2 \omega d}{c}} E_0^{IV}$

de 1,2

$$\begin{aligned}
 \sqrt{\epsilon_1} (E_0 + E_0^{IV}) &= \sqrt{\epsilon_2} (E_0^I + E_0^{III} - E_0 + E_0^{IV}) \\
 E_0^{IV} (\sqrt{\epsilon_1} - \sqrt{\epsilon_2}) &= \sqrt{\epsilon_2} 2 E_0^I - (\sqrt{\epsilon_2} + \sqrt{\epsilon_1}) E_0 \\
 E_0^{IV} &= \frac{2 E_0^I \sqrt{\epsilon_2} - E_0 (\sqrt{\epsilon_2} + \sqrt{\epsilon_1})}{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}}
 \end{aligned}$$

de 3,4

$$\begin{aligned}
 \sqrt{\epsilon_2} \cdot E_0 + \sqrt{\epsilon_2} \cdot E_0^{IV} \cdot e^{-i \frac{2 \omega d}{c} \sqrt{\epsilon_2}} &= \sqrt{\epsilon_3} \cdot E_0^I - \sqrt{\epsilon_3} \cdot E_0^{III} \cdot e^{i \frac{2 \omega d}{c} \sqrt{\epsilon_3}} \\
 E_0^{IV} \cdot e^{-i \frac{2 \omega d}{c} \sqrt{\epsilon_2}} (\sqrt{\epsilon_2} + \sqrt{\epsilon_3}) &= E_0^I (\sqrt{\epsilon_3} - \sqrt{\epsilon_2}) \\
 (E_0^{IV} - E_0 + E_0^I) e^{-i \frac{2 \omega d}{c} \sqrt{\epsilon_2}} &= E_0^I (\sqrt{\epsilon_3} - \sqrt{\epsilon_2}) \\
 E_0^{IV} &= E_0^I \frac{(\sqrt{\epsilon_3} - \sqrt{\epsilon_2})}{(\sqrt{\epsilon_2} + \sqrt{\epsilon_3})} \cdot e^{i \frac{2 \omega d}{c} \sqrt{\epsilon_2}} - E_0 + E_0 \\
 \frac{2 E_0^I \sqrt{\epsilon_2} - E_0 (\sqrt{\epsilon_2} + \sqrt{\epsilon_1})}{(\sqrt{\epsilon_1} - \sqrt{\epsilon_2})} &= E_0^I \left(\frac{(\sqrt{\epsilon_3} - \sqrt{\epsilon_2})}{(\sqrt{\epsilon_2} + \sqrt{\epsilon_3})} e^{i \frac{2 \omega d}{c} \sqrt{\epsilon_2}} - 1 \right) + E_0 \\
 E_0^I \left(\frac{2 \sqrt{\epsilon_2}}{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}} - \frac{\sqrt{\epsilon_3} - \sqrt{\epsilon_2}}{(\sqrt{\epsilon_2} + \sqrt{\epsilon_3})} e^{i \frac{2 \omega d}{c} \sqrt{\epsilon_2}} + 1 \right) &= E_0 \left(1 + \frac{\sqrt{\epsilon_2} + \sqrt{\epsilon_1}}{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}} \right) = \frac{E_0 2 \sqrt{\epsilon_1}}{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}} \\
 E_0^I &= \frac{2 E_0 \sqrt{\epsilon_1}}{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}} \cdot \frac{1}{\left[\frac{2 \sqrt{\epsilon_2}}{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}} - \frac{\sqrt{\epsilon_3} - \sqrt{\epsilon_2}}{(\sqrt{\epsilon_2} + \sqrt{\epsilon_3})} e^{i \frac{2 \omega d}{c} \sqrt{\epsilon_2}} + 1 \right]} \\
 E_0^{IV} &= \frac{\frac{2 \sqrt{\epsilon_2}}{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}} \cdot 2 E_0 \sqrt{\epsilon_1} \cdot \frac{1}{\left[\frac{2 \sqrt{\epsilon_2}}{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}} - \frac{\sqrt{\epsilon_3} - \sqrt{\epsilon_2}}{(\sqrt{\epsilon_2} + \sqrt{\epsilon_3})} e^{i \frac{2 \omega d}{c} \sqrt{\epsilon_2}} + 1 \right]} - E_0 (\sqrt{\epsilon_2} + \sqrt{\epsilon_1})}{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}} = 0 \\
 \frac{4 \sqrt{\epsilon_2} \sqrt{\epsilon_1}}{(\sqrt{\epsilon_1} - \sqrt{\epsilon_2})} &= \sqrt{\epsilon_2} + \sqrt{\epsilon_1} \rightarrow \\
 \frac{4 \sqrt{\epsilon_1} \sqrt{\epsilon_2}}{(\sqrt{\epsilon_1} - \sqrt{\epsilon_2}) (\sqrt{\epsilon_1} + \sqrt{\epsilon_2})} &= \frac{1}{\sqrt{\epsilon_1} + \sqrt{\epsilon_2}}
 \end{aligned}$$

$$\frac{\sqrt{E_1 E_2}}{(\sqrt{E_1} - \sqrt{E_2})(\sqrt{E_1} + \sqrt{E_2})} = \frac{Z\sqrt{E_2}}{(\sqrt{E_1} - \sqrt{E_2})} e^{i \frac{Z\omega d \sqrt{E_2}}{c}} + 1$$

$$\frac{2\sqrt{E_2}}{\sqrt{E_1} - \sqrt{E_2}} \left(\frac{\sqrt{E_1} Z}{\sqrt{E_1} + \sqrt{E_2}} - 1 \right) = \dots e + 1$$

$$\frac{2\sqrt{E_2}}{\sqrt{E_1} + \sqrt{E_2}} \left(\frac{Z\sqrt{E_1}}{\sqrt{E_1} + \sqrt{E_2}} - 1 \right) = 1 - \frac{(\sqrt{E_3} - \sqrt{E_2})}{(\sqrt{E_2} + \sqrt{E_3})} e^{i \frac{Z\omega d \sqrt{E_2}}{c}}$$

$$\frac{2\sqrt{E_2}}{\sqrt{E_1} - \sqrt{E_2}} \left(\frac{Z\sqrt{E_1} - \sqrt{E_1} - \sqrt{E_2}}{\sqrt{E_1} + \sqrt{E_2}} \right) = \frac{2\sqrt{E_2}}{\sqrt{E_1} + \sqrt{E_2}} =$$

$$\frac{2\sqrt{E_2} - \sqrt{E_1} - \sqrt{E_2}}{\sqrt{E_1} + \sqrt{E_2}} = \frac{\sqrt{E_2} - \sqrt{E_1}}{(\sqrt{E_2} + \sqrt{E_3})} e^{i \frac{Z\omega d \sqrt{E_2}}{c}}$$

$$\frac{(\sqrt{E_2} - \sqrt{E_1})}{\sqrt{E_1} + \sqrt{E_2}} \frac{(\sqrt{E_2} + \sqrt{E_3})}{(\sqrt{E_2} - \sqrt{E_3})} = e^{i \frac{Z\omega d \sqrt{E_2}}{c}}$$

$$= \begin{cases} \cos\left(\frac{Z\omega d \sqrt{E_2}}{c}\right) \\ i \cdot \sin\left(\frac{Z\omega d \sqrt{E_2}}{c}\right) \end{cases}$$

$$\frac{\sqrt{E_2} - \sqrt{E_1}}{(\sqrt{E_1} + \sqrt{E_2})} \frac{\sqrt{E_2} + \sqrt{E_3}}{(\sqrt{E_2} - \sqrt{E_3})} = 1$$

* 1, m par

$$E_2 - \sqrt{E_1} E_2 + \sqrt{E_2} E_3 - \sqrt{E_1} E_3 =$$

$$\sqrt{E_1} E_2 + \sqrt{E_1} E_3 - \sqrt{E_2} E_3$$

$$\Rightarrow -\sqrt{E_1} E_2 + \sqrt{E_2} E_3 = -\sqrt{E_2} E_3 + \sqrt{E_1} E_2$$

$$\cancel{\sqrt{E_2} E_3} = \cancel{\sqrt{E_1} E_2}$$

$$\boxed{E_3 = E_1}$$

* -1, m impar

$$E_2 - \sqrt{E_1} E_2 + \sqrt{E_2} E_3 - \sqrt{E_1} E_3 =$$

$$-\sqrt{E_1} E_2 - E_2 + \sqrt{E_1} E_3 + \sqrt{E_2} E_3$$

$$\cancel{\sqrt{E_2} E_3} = \cancel{\sqrt{E_1} E_3}$$

$$\boxed{E_2 = \sqrt{E_1} E_3}$$

$$\frac{Z\omega d \sqrt{E_2}}{c} = m\pi$$

$$\boxed{d = \frac{m \pi \cdot c}{Z\omega \sqrt{E_2}}} \quad m \in \mathbb{Z}$$

$$\cancel{E_2 \sqrt{E_2}} + \cancel{E_1 \sqrt{E_2}} + \cancel{E_2 \sqrt{E_3}} + \cancel{E_1 \sqrt{E_3}} =$$

$$- \cancel{E_2 \sqrt{E_3}} + \cancel{E_1 \sqrt{E_3}} + \cancel{E_2 \sqrt{E_2}} - \cancel{E_1 \sqrt{E_2}}$$

$$Z E_2 \sqrt{E_3} = -Z E_1 \sqrt{E_2}$$

$$\frac{\sqrt{E_3}}{E_1} = -\frac{1}{\sqrt{E_2}}$$

$$E_2 = -\frac{E_1^2}{E_3}$$

4.

$$\vec{E} = \vec{E}_0 \cdot e^{i(\vec{k} \cdot \vec{x} - \omega t)}$$

$$\vec{H} = \frac{\vec{B}}{\mu} = \frac{c}{\mu \omega} \vec{k} \times \vec{E}$$

$$\mu_{\text{em}} = \frac{1}{8\pi} (\vec{E} \cdot \vec{D} + \vec{B} \cdot \vec{H}) \Rightarrow$$

MLIH (μ, ϵ)
homogéneas

$$\langle \mu_e \rangle = \frac{1}{8\pi} \text{Re} \left(\frac{\vec{E} \cdot \vec{D}^*}{2} \right)$$

$$\langle \mu_m \rangle = \frac{1}{8\pi} \text{Re} \left(\frac{\vec{B} \cdot \vec{H}^*}{2} \right)$$

Pero, en un conductor $\vec{k} \in \mathbb{C}$ y es

$$\vec{k} = k \sqrt{1 + \frac{c \cdot 4\pi\sigma}{\omega \epsilon}} = \sqrt{\mu \epsilon} \cdot \frac{\omega}{c} \sqrt{1 + i \frac{4\pi\sigma}{\omega \epsilon}}$$

$$\vec{k} = \left(\beta + i \frac{\alpha}{2} \right) \hat{k} \quad \text{suponemos propagación en } x$$

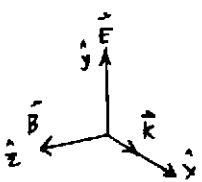
$$\vec{E} = \vec{E}_0 \cdot e^{i\left[\left(\beta + i \frac{\alpha}{2}\right) \cdot x - \omega t\right]} = \vec{E}_0 \cdot e^{-\frac{\alpha}{2}x} \cdot e^{i(\beta x - \omega t)}$$

$$\text{con } \frac{\alpha}{2} = \sqrt{\mu \epsilon} \cdot \frac{\omega}{c} \cdot \left(\frac{\sqrt{1 + \left(\frac{4\pi\sigma}{\omega \epsilon}\right)^2} - 1}{2} \right)^{1/2}$$

$$\langle \mu_e \rangle = \frac{1}{8\pi} \cdot \frac{\vec{E}_0}{2} \cdot e^{-\frac{\alpha}{2}x} \cdot e^{i(\beta x - \omega t)} \cdot \epsilon \cdot \vec{E}_0 \cdot e^{-\frac{\alpha}{2}x} \cdot e^{-i(\beta x - \omega t)}$$

$$\langle \mu_e \rangle = \frac{1}{16\pi} |\vec{E}_0|^2 \cdot e^{-\alpha x} \cdot \epsilon \quad \leftarrow \text{energía asociada al campo eléctrico}$$

$$\vec{H} = \frac{c}{\mu \omega} \vec{k} \times \vec{E} = \frac{c}{\mu \omega} \left(\beta + i \frac{\alpha}{2} \right) \cdot \vec{E}_0 \cdot e^{-\frac{\alpha}{2}x} \cdot e^{i(\beta x - \omega t)}$$



$$\langle \mu_m \rangle = \frac{1}{8\pi} \cdot \frac{c}{2\omega} \vec{E}_0 \cdot e^{-\frac{\alpha}{2}x} \cdot \frac{c}{\mu \omega} \vec{E}_0 \cdot e^{-\frac{\alpha}{2}x} \cdot \left(\beta^2 + \frac{\alpha^2}{4} \right)$$

forma de \vec{H}

$$\frac{c}{\mu \omega} \vec{E}_0 \cdot e^{-\frac{\alpha}{2}x} \cdot \left(\beta + i \frac{\alpha}{2} \right) e^{i(\beta x - \omega t)}$$

$$\cdot e^{i(\beta x - \omega t)} \cdot \left(\beta^2 + \frac{\alpha^2}{4} \right)^{1/2} \cdot e^{i \tan^{-1} \left(\frac{\alpha}{2\beta} \right)}$$

$$\frac{c}{\mu \omega} \vec{E}_0 \cdot e^{-\frac{\alpha}{2}x} \cdot e^{i\left[\beta x - \omega t + \tan^{-1} \left(\frac{\alpha}{2\beta} \right)\right]} \cdot \left(\beta^2 + \frac{\alpha^2}{4} \right)^{1/2}$$

Relación entre las energías asociadas a los campos

$$\langle \mu_m \rangle = \frac{1}{16\pi} \frac{c^2}{\mu \omega^2} |\vec{E}_0|^2 \cdot e^{-\alpha x} \cdot \left(\beta^2 + \frac{\alpha^2}{4} \right)$$

$$\langle \mu_m \rangle = \frac{c^2}{\mu \omega^2} \cdot \frac{\langle \mu_e \rangle}{\epsilon} \cdot \left(\beta^2 + \frac{\alpha^2}{4} \right)$$

$$\langle \mu_m \rangle = \frac{c^2}{\omega^2} \cdot \frac{1}{\mu \epsilon} \langle \mu_e \rangle \cdot \left(\beta^2 + \frac{\alpha^2}{4} \right)$$

$$\langle \mu_m \rangle = \frac{1}{k^2} \cdot \left(\beta^2 + \frac{\alpha^2}{4} \right) \langle \mu_e \rangle$$

a. Supongamos un mal conductor (dieléctrico)

$$\Rightarrow \frac{4\pi\sigma}{\omega\epsilon} \ll 1 \rightarrow \left(1 + \left(\frac{4\pi\sigma}{\omega\epsilon}\right)^2\right)^{1/2} = 1 + \frac{1}{2} \left(\frac{4\pi\sigma}{\omega\epsilon}\right)^2$$

$$\beta = \sqrt{\mu\epsilon} \cdot \frac{\omega}{c} \cdot \left(\frac{\sqrt{1 + \left(\frac{4\pi\sigma}{\omega\epsilon}\right)^2} + 1}{2}\right)^{1/2} \quad \frac{\alpha}{Z} = \sqrt{\mu\epsilon} \cdot \frac{\omega}{c} \cdot \left(\frac{\sqrt{1 + \left(\frac{4\pi\sigma}{\omega\epsilon}\right)^2} - 1}{2}\right)^{1/2}$$

$$\left(\frac{1 + \frac{1}{2} \left(\frac{4\pi\sigma}{\omega\epsilon}\right)^2 + 1}{2}\right)^{1/2} \quad \left(\frac{1 + \frac{1}{2} \left(\frac{4\pi\sigma}{\omega\epsilon}\right)^2 - 1}{2}\right)^{1/2}$$

$$\left[1 + \frac{1}{4} \left(\frac{4\pi\sigma}{\omega\epsilon}\right)^2\right]^{1/2} \quad \left[\frac{1}{4} \left(\frac{4\pi\sigma}{\omega\epsilon}\right)^2\right]^{1/2}$$

$$\beta \approx \sqrt{\mu\epsilon} \cdot \frac{\omega}{c}$$

$$\beta^2 \approx \mu\epsilon \frac{\omega^2}{c^2}$$

$$\frac{\alpha}{Z} \approx \sqrt{\mu\epsilon} \cdot \frac{\omega}{c} \cdot \frac{1}{2} \cdot \frac{4\pi\sigma}{\omega\epsilon}$$

$$\frac{\alpha}{Z} \approx \frac{4\pi\sigma\mu}{\sqrt{\epsilon} \cdot 2c}$$

$$\frac{\alpha^2}{4} \approx \frac{4\pi^2 \sigma^2 \mu^2}{\epsilon c^2}$$

$$\left(\beta^2 + \frac{\alpha^2}{4}\right) = \mu\epsilon \frac{\omega^2}{c^2} + \frac{4\pi^2 \sigma^2 \mu^2}{\epsilon c^2}$$

$$\langle \mu_M \rangle = \frac{1}{\mu\epsilon \frac{\omega^2}{c^2}} \left(\mu\epsilon \frac{\omega^2}{c^2} + \frac{4\pi^2 \sigma^2 \mu^2}{\epsilon c^2} \right) \langle \mu_e \rangle = \left(1 + \frac{4\pi^2 \sigma^2}{\omega^2 \epsilon^2} \right) \langle \mu_e \rangle$$

$$\langle \mu_M \rangle = \left(1 + \frac{1}{4} \left[\frac{4\pi\sigma}{\omega\epsilon} \right]^2 \right) \langle \mu_e \rangle$$

b. Supongamos ahora un buen conductor (muy buen conductor)

$$\frac{4\pi\sigma}{\omega\epsilon} \gg 1 \Rightarrow 1 + \left(\frac{4\pi\sigma}{\omega\epsilon}\right)^2 \approx \left(\frac{4\pi\sigma}{\omega\epsilon}\right)^2$$

$$\beta \approx \sqrt{\mu\epsilon} \cdot \frac{\omega}{c} \cdot \left(\frac{4\pi\sigma}{\omega\epsilon}\right)^{1/2}$$

$$\frac{\alpha}{Z} \approx \sqrt{\mu\epsilon} \cdot \frac{\omega}{c} \cdot \left(\frac{4\pi\sigma}{\omega\epsilon}\right)^{1/2}$$

$$\sqrt{\mu\epsilon} \cdot \frac{\omega}{c} \left(\frac{4\pi\sigma}{\omega\epsilon}\right)^{1/2}$$

$$\beta^2 + \frac{\alpha^2}{4} = \left(\mu\epsilon \cdot \frac{\omega^2}{c^2} \cdot \frac{4\pi\sigma}{\omega\epsilon} \right)^2 \Rightarrow$$

$$\langle \mu_M \rangle = \frac{1}{\cancel{\mu\epsilon} \frac{\omega^2}{c^2}} \cdot \cancel{\mu\epsilon} \frac{\omega^2}{c^2} \cdot \frac{4\pi\sigma}{\omega\epsilon} \cdot 2 \langle \mu_e \rangle = \frac{4\pi\sigma}{\omega\epsilon} \langle \mu_e \rangle$$

$$\langle \mu_M \rangle = \frac{4\pi\sigma}{\omega\epsilon} \langle \mu_e \rangle$$

5.

$\vec{A}(\vec{x}, t), \vec{B}(\vec{x}, t)$ armónicos \Rightarrow

$$\vec{A}(\vec{x}, t) = \vec{A}(\vec{x}) \cdot e^{i\omega t}$$

$$\vec{B}(\vec{x}, t) = \vec{B}(\vec{x}) \cdot e^{i\omega t}$$

} con $\vec{A}(\vec{x}), \vec{B}(\vec{x}) \in \mathbb{C}$

$$\text{Re}\{\vec{A}(\vec{x}, t)\} \cdot \text{Re}\{\vec{B}(\vec{x}, t)\} =$$

$$|\vec{A}(\vec{x})| \cdot \cos(\omega t + \varphi_A) \cdot \cos(\omega t + \varphi_B) \cdot |\vec{B}(\vec{x})|$$

→ Pero esto ya es muy complicado por las fases de cada complejo

$$\text{Re}\{\vec{A}(\vec{x}, t)\} = \frac{\vec{A}(\vec{x}) \cdot e^{i\omega t} + \vec{A}^*(\vec{x}) \cdot e^{-i\omega t}}{2}$$

$$\text{Re}\{\vec{B}(\vec{x}, t)\} = \frac{\vec{B}(\vec{x}) \cdot e^{i\omega t} + \vec{B}^*(\vec{x}) \cdot e^{-i\omega t}}{2}$$

$$\Rightarrow \text{Re}\{\vec{A}(\vec{x}, t)\} \cdot \text{Re}\{\vec{B}(\vec{x}, t)\} =$$

$$\frac{1}{4} \left(\vec{A}(\vec{x}) \cdot \vec{B}(\vec{x}) \cdot e^{2i\omega t} + \vec{A}(\vec{x}) \cdot \vec{B}^*(\vec{x}) + \vec{A}^*(\vec{x}) \cdot \vec{B}(\vec{x}) + \vec{A}^*(\vec{x}) \cdot \vec{B}^*(\vec{x}) \cdot e^{-2i\omega t} \right) =$$

$$\frac{1}{4} \left(2 \text{Re}\{\vec{A}(\vec{x}) \cdot \vec{B}(\vec{x}) \cdot e^{2i\omega t}\} + 2 \text{Re}\{\vec{A}(\vec{x}) \cdot \vec{B}^*(\vec{x})\} \right) =$$

$$\frac{1}{2} \text{Re}\{\vec{A}(\vec{x}) \cdot \vec{B}^*(\vec{x}) + \vec{A}(\vec{x}) \cdot \vec{B}(\vec{x}) \cdot e^{2i\omega t}\} \Rightarrow$$

$$\langle \text{Re } \vec{A} \cdot \text{Re } \vec{B} \rangle = \frac{1}{2} \text{Re}\{\vec{A}(\vec{x}) \cdot \vec{B}^*(\vec{x}) + \vec{A}(\vec{x}) \cdot \vec{B}(\vec{x}) \cdot \int_0^{2\pi} e^{2i\omega t} dt \cdot \frac{1}{n}\}$$

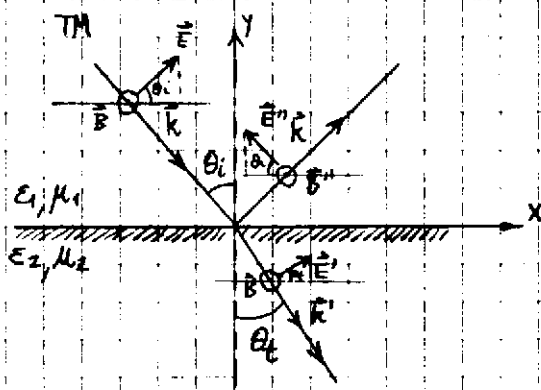
integral

$$\begin{aligned} & \begin{matrix} z i \omega t = 0 \\ z i \omega t = 2\pi n \\ dt = -\frac{du}{2\omega} \end{matrix} \rightarrow -\int_0^{2\pi n} \frac{e^u}{2\omega} \cdot \frac{1}{n} = \\ & -\left(e^{2\pi n} - e^0 \right) \cdot \frac{1}{n} = \\ & \left[1 - (\cos(2\pi n) + i \sin 2\pi n) \right] \cdot \frac{1}{n} = \\ & [1 - 1] \cdot \frac{1}{n} = 0 \end{aligned}$$

$$\begin{aligned} z i \omega t &= 2\pi n \\ t &= \frac{2\pi n}{\omega} \end{aligned}$$

$$\langle \text{Re}(\vec{A}) \cdot \text{Re}(\vec{B}) \rangle = \frac{1}{2} \text{Re}(\vec{A}(\vec{x}) \cdot \vec{B}^*(\vec{x}))$$

$$P_{rad} = \frac{dF_{mec}}{dS}$$



$$\frac{d\vec{P}}{dt} = \oint_S \vec{T} \cdot d\vec{S}$$

$$\text{con } \vec{P} = \vec{P}_{mec} + \vec{P}_{campo}$$

$$\vec{P}_{mec} = \oint_S \vec{T} \cdot d\vec{S} - \frac{d}{dt} \int_V \frac{\vec{E} \times \vec{B}}{4\pi c} dV$$

La superficie S debe usar la normal \hat{n} ; pero dado nuestro sistema de coordenadas invertiremos el signo y será $\frac{d\vec{P}}{dt} = - \oint_S \vec{T} \cdot d\vec{S}$ con $\hat{n} \equiv \hat{y}$

Para la presión me interesará la fuerza en \hat{y}

* condiciones

$$\vec{D}_n \quad [\epsilon_1 (\vec{E}_0 + \vec{E}_0'') - \epsilon_2 \vec{E}_0'] \cdot \hat{n} = 0$$

$$\vec{B}_t \quad \frac{1}{\mu_1} (k E_0 + k E_0'') - \frac{1}{\mu_2} k' E_0' = 0$$

$$\epsilon_1 \sin(\theta_i) (E_0 + E_0'') = \epsilon_2 E_0' \sin(\theta_t)$$

$$\sqrt{\frac{\epsilon_1}{\mu_1}} (E_0 + E_0'') = \sqrt{\frac{\epsilon_2}{\mu_2}} E_0'$$

$$\vec{E}_t \quad [(\vec{E}_0 + \vec{E}_0'') - \vec{E}_0'] \times \hat{n} = 0$$

$$\cos(\theta_i) E_0 - \cos(\theta_r) E_0'' = \cos(\theta_t) E_0'$$

$$\vec{B} = \sqrt{\mu} \epsilon \frac{\vec{k} \times \vec{E}}{k}$$

$$B_0 = \sqrt{\mu} \epsilon E_0$$

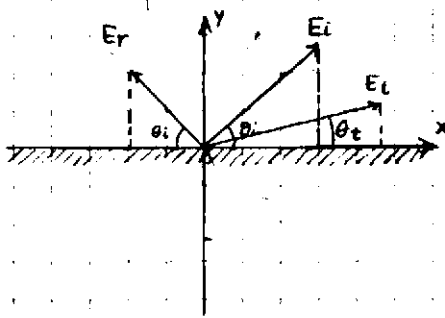
definimos

$$\begin{aligned} E_0 &\equiv E_i \\ E_0'' &\equiv E_r \\ E_0' &\equiv E_t \end{aligned}$$

$$\begin{cases} \epsilon_1 \sin \theta_i (E_i + E_r) = \epsilon_2 E_t \sin \theta_t \\ \cos \theta_i (E_i - E_r) = \cos \theta_t E_t \end{cases}$$

$$\begin{aligned} B_i &= \sqrt{\mu_1 \epsilon_1} E_i \\ B_r &= \sqrt{\mu_1 \epsilon_1} E_r \\ B_t &= \sqrt{\mu_2 \epsilon_2} E_t \end{aligned}$$

En el plano ZX



$$\vec{E}_r \cdot \hat{y} = |E_r| \sin \theta_i$$

$$\vec{E}_i \cdot \hat{y} = |E_i| \sin \theta_i$$

$$\vec{E}_t \cdot \hat{y} = |E_t| \sin \theta_t$$

en el plano ZX

es:

$$\begin{aligned} \vec{E} &= E \hat{y} + E \hat{x} \\ \vec{B} &= B \hat{z} \end{aligned}$$

$$\vec{T} = \begin{pmatrix} T_{xx} & T_{xy} & T_{xz} \\ T_{yx} & T_{yy} & T_{yz} \\ T_{zx} & T_{zy} & T_{zz} \end{pmatrix}$$

$$d\vec{S} = \begin{pmatrix} 0 \\ dx dz \\ 0 \end{pmatrix}$$

$$d\vec{F} \cdot \hat{y} = T_{yy} dS$$

$$d\vec{F} \cdot \hat{y} = \frac{1}{4\pi} \left(\epsilon \frac{E_y^2}{2} - \frac{\epsilon E_x^2}{2} - \frac{\epsilon E_z^2}{2} - \frac{\mu}{2} \frac{B_z^2}{\mu x} \right) dS \hat{y}$$

$$B_z^2 = (\sqrt{\mu_1 \epsilon_1} (E_i + E_r))^2$$

la Normal es en \hat{y}

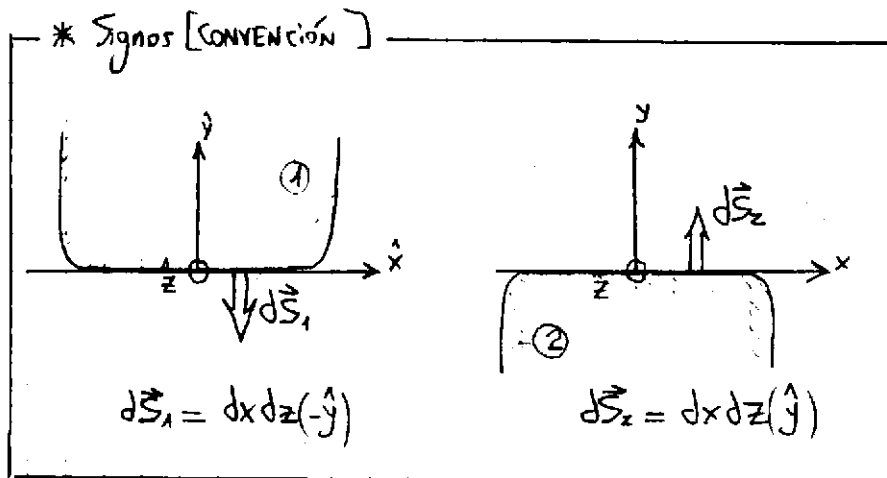
$$\begin{aligned} & - \frac{1}{4\pi} \left[\frac{\epsilon_1}{2} \sin^2 \theta_i (E_i + E_r)^2 + \frac{\epsilon_2}{2} (E_i - E_r)^2 \cos^2 \theta_i - \frac{1}{2\mu_1} (E_i + E_r)^2 \mu_1 \epsilon_1 \right] \\ & + \frac{1}{4\pi} \left[\frac{\epsilon_2}{2} \sin^2 \theta_t E_t^2 - \frac{\epsilon_2}{2} E_t^2 \cos^2 \theta_t - \frac{\mu_2}{2} \mu_2 \epsilon_2 E_t^2 \right] \end{aligned}$$

$$= \frac{1}{8\pi} \left(-\epsilon_1 \sin^2 \theta_i (E_i + E_r)^2 + \epsilon_1 (E_i - E_r)^2 \cos^2 \theta_i + \epsilon_1 (E_i + E_r)^2 - \epsilon_2 \sin^2 \theta_t E_t^2 + \epsilon_2 E_t^2 \cos^2 \theta_t + \epsilon_2 E_t^2 \right)$$

$$= \frac{1}{8\pi} \left[-(E_i + E_r)^2 \epsilon_1 [1 - \sin^2 \theta_i] + \epsilon_1 (E_i - E_r)^2 \cos^2 \theta_i - \epsilon_2 E_t^2 (1 + \cos^2 \theta_t - \sin^2 \theta_t) \right]$$

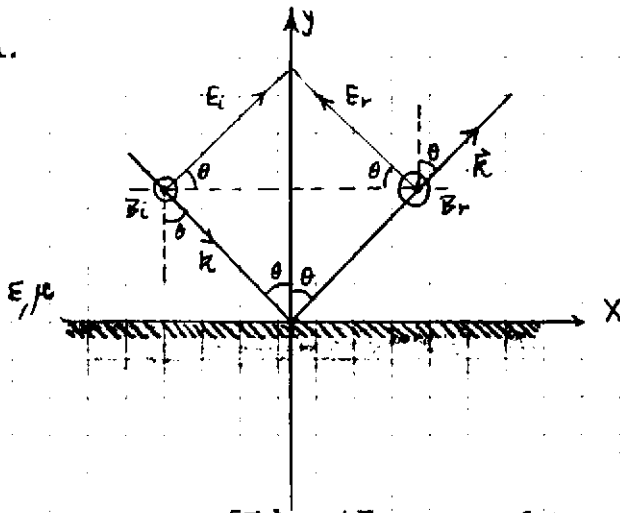
$$d\vec{F} \cdot \hat{y} = \frac{1}{8\pi} \left[2\epsilon_1 \cos^2 \theta_i (E_i^2 + E_r^2) + \epsilon_2 E_t^2 \cdot 2 \cos^2 \theta_t \right]$$

$$\text{Presión Radiación} = \frac{1}{4\pi} \left([|E_i|^2 + |E_r|^2] \epsilon_1 \cos^2 \theta_i + |E_t|^2 \epsilon_2 \cos^2 \theta_t \right)$$



6.

a.



Onda plana incide sobre espejo perfecto con ángulo $\theta \Rightarrow$ no hay onda transmitida

$$k_{inc} = k_{refl}$$

$$\vec{E}_i = E_i e^{i[\vec{k} \cdot \vec{x} - \omega t]} = E_i e^{i[-k \cdot y \cos \theta + k \cdot x \sin \theta - \omega t]} (\cos \theta \hat{x} + \sin \theta \hat{y})$$

$$\vec{E}_r = E_r e^{i[k \cdot y \cos \theta + k \cdot x \sin \theta - \omega t]} (-\cos \theta \hat{x} + \sin \theta \hat{y})$$

$$\vec{B}_i = \sqrt{\mu \epsilon} \hat{k} \times \vec{E}_i = \sqrt{\mu \epsilon} E_i e^{i[-k \cdot y \cos \theta + k \cdot x \sin \theta - \omega t]} \hat{z}$$

$$\vec{B}_r = \sqrt{\mu \epsilon} E_r e^{i[k \cdot y \cos \theta + k \cdot x \sin \theta - \omega t]} \hat{z}$$

* Continuidad

\vec{E}_t, \vec{B}_n continuos \rightarrow

$\nabla \cdot \vec{H}, \nabla \times \vec{E}$ no pueden usarse porque hay \vec{j}, σ superficiales

$$E_i \cos \theta - E_r \cos \theta = 0$$

$$E_i = E_r$$

* Tensor de Maxwell

$$dS_j = d\vec{S} = \begin{pmatrix} 0 \\ -dx \cdot dz \\ 0 \end{pmatrix} \Rightarrow \vec{T} \cdot d\vec{S} = \begin{pmatrix} T_{xx} & T_{xy} & T_{xz} \\ T_{yx} & T_{yy} & T_{yz} \\ T_{zx} & T_{zy} & T_{zz} \end{pmatrix} \cdot \begin{pmatrix} 0 \\ -dS \\ 0 \end{pmatrix} = \begin{pmatrix} -T_{xy} \cdot dS \\ -T_{yy} \cdot dS \\ -T_{zy} \cdot dS \end{pmatrix}$$

$$-T_{yy} \cdot dS = -\frac{1}{4\pi} \left(\underbrace{\epsilon E_y^2}_{=0} + \underbrace{\mu H_y^2}_{=0} - \frac{\epsilon}{2} (E_x^2 + E_y^2 + E_z^2) - \frac{\mu}{2} (H_x^2 + H_y^2 + H_z^2) \right) \cdot dS$$

$$-T_{yy} \cdot dS = -\frac{1}{4\pi} \left(\frac{\epsilon}{2} E_y^2 - \frac{\epsilon}{2} E_x^2 - \frac{\mu}{2} H_z^2 \right) dS \Rightarrow$$

$$\langle -T_{yy} \rangle = \left(\frac{\epsilon}{4} [\text{Re}(E_y \cdot E_y^*) - \text{Re}(E_x \cdot E_x^*)] - \frac{\mu}{4} [\text{Re}(H_z \cdot H_z^*)] \right) \cdot \left(-\frac{1}{4\pi} \right)$$

$$E_y \cdot E_y^* = \left[E_i e^{i(ky \cos \theta + kx \sin \theta - \omega t)} \cdot \text{sen} \theta + E_i e^{i(ky \cos \theta - kx \sin \theta + \omega t)} \cdot \text{sen} \theta \right] \cdot \left[E_i e^{i(ky \cos \theta + kx \sin \theta - \omega t)} \cdot \text{sen} \theta + E_i e^{i(ky \cos \theta - kx \sin \theta + \omega t)} \cdot \text{sen} \theta \right]$$

$$= E_i \cdot \text{sen} \theta \cdot e^{i(kx \sin \theta - \omega t)} \cdot \left(e^{-iky \cos \theta} + e^{iky \cos \theta} \right) + E_i \cdot \text{sen} \theta \cdot e^{-i(kx \sin \theta - \omega t)} \cdot \left(e^{iky \cos \theta} + e^{-iky \cos \theta} \right)$$

$$= E_i^2 \cdot \text{sen}^2 \theta \cdot (\cos[k \cdot y \cdot \cos \theta]) \cdot (\cos[k \cdot y \cdot \cos \theta]) \cdot 4 \Rightarrow E_y \cdot E_y^* \Big|_{y=0} = 4 E_i^2 \text{sen}^2 \theta$$

$$E_x \cdot E_x^* = E_i \cos \theta \cdot e^{i(kx \sin \theta - \omega t)} \left[e^{-iky \cos \theta} - e^{iky \cos \theta} \right] \\ E_i \cos \theta \cdot e^{-i(kx \sin \theta - \omega t)} \left[e^{iky \cos \theta} - e^{-iky \cos \theta} \right] = 0 = E_x \cdot E_x^* \Big|_{y=0}$$

$$H_z \cdot H_z^* = \frac{B_z B_z^*}{\mu^2} = \frac{\epsilon \cdot E_i^2}{\mu^2}$$

→

$$\langle -T_{yy} \rangle = \left(\frac{\epsilon}{4} \cdot 4 E_i^2 \sin^2 \theta - \frac{\epsilon}{4} \cdot E_i^2 \right) \cdot \frac{1}{4\pi} \\ - \frac{\epsilon E_i^2}{4} \left[4 \sin^2 \theta - 1 \right] \cdot \frac{1}{4\pi} = \frac{\epsilon}{16\pi} E_i^2 (1 - 4 \sin^2 \theta)$$

$$\text{Presión Radiación} = \langle -T_{yy} \rangle = \frac{\epsilon}{16\pi} E_i^2 (1 - 4 \sin^2 \theta)$$

b. Si $\theta = 0$ (incidencia normal) ⇒

$$\text{Presión Radiación} = \frac{\epsilon}{16\pi} E_i^2$$

$$\langle \mu_{em} \rangle = \frac{1}{8\pi} \cdot \frac{1}{2} \text{Re} (E \cdot D^* + B \cdot H^*)$$

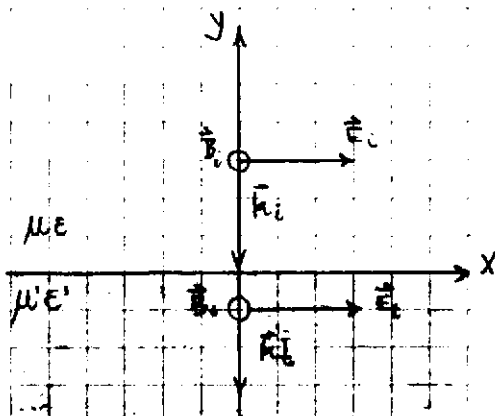
$$\frac{1}{16\pi} \text{Re} \left\{ \begin{aligned} & \epsilon E_i^2 \sin^2 \theta \cdot e^{i(kx \sin \theta - \omega t)} \cdot 2 \cos^2(ky \cos \theta) \cdot e^{-i(kx \sin \theta - \omega t)} \\ & \epsilon E_i^2 \cos^2 \theta \cdot e^{i(kx \sin \theta - \omega t)} \cdot 2 \sin^2(ky \cos \theta) \cdot e^{-i(kx \sin \theta - \omega t)} \\ & + \cancel{\mu E_i^2} \cdot e^{i(\dots)} \cdot e^{-i(\dots)} \end{aligned} \right\}$$

$$\langle \mu_{em} \rangle = \frac{1}{16\pi} \left(\epsilon E_i^2 \sin^2 \theta \cdot 4 \cos^2(ky \cos \theta) - \epsilon E_i^2 \cos^2 \theta \cdot 4 \sin^2(ky \cos \theta) + \epsilon E_i^2 \right)$$

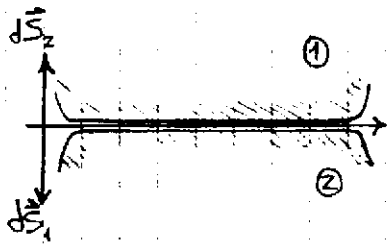
Pero si $\theta = 0 \rightarrow$

$$\mu_{em} = \frac{1}{16\pi} \epsilon \cdot E_i^2$$

c.



$$\vec{E}_i = E_i \cdot e^{i(k \cdot y - \omega t)} \hat{x} \\ \vec{B}_i = \sqrt{\mu \epsilon} E_i \cdot e^{i(k \cdot y - \omega t)} \hat{z} \\ \vec{E}_t = E_t \cdot e^{i(-k \cdot y - \omega t)} \hat{x} \\ \vec{B}_t = \sqrt{\mu' \epsilon'} E_t \cdot e^{i(-k \cdot y - \omega t)} \hat{z}$$



$$d\vec{F} = \vec{T} \cdot d\vec{S} = \begin{pmatrix} T_{xx} & T_{yx} & T_{yz} \\ T_{yx} & T_{yy} & T_{yz} \\ T_{yz} & T_{yz} & T_{zz} \end{pmatrix} \begin{pmatrix} 0 \\ -dx dz \\ 0 \end{pmatrix} = \begin{pmatrix} -T_{xy} \cdot dx dz \\ -T_{yy} \cdot dx dz \\ -T_{zy} \cdot dx dz \end{pmatrix}$$

Presión = $\langle -T_{yy} \rangle$

* medio ①

$$(4\pi) \cdot -T_{yy} = - \left[\epsilon E_y^2 + \mu H_y^2 - \frac{\epsilon}{2} (E_y^2 + E_z^2 + E_x^2) - \frac{\mu}{2} (H_x^2 + H_y^2 + H_z^2) \right]$$

$$(4\pi) \cdot -T_{yy} = \left[-\frac{\epsilon E_y^2}{2} + \frac{\epsilon E_x^2}{2} + \frac{\epsilon E_z^2}{2} - \frac{\mu H_y^2}{2} + \frac{\mu H_x^2}{2} + \frac{\mu H_z^2}{2} \right]$$

Pero $\vec{E} = E_x \hat{x}$, $\vec{B} = B_z \hat{z} \rightarrow -T_{yy} = \left(\frac{\epsilon E_x^2}{2} + \frac{\mu H_z^2}{2} \right) \frac{1}{4\pi} \therefore$

$$\langle -T_{yy} \rangle = \frac{1}{8\pi} \cdot \frac{1}{2} \text{Re} \left\{ \epsilon E_x \cdot E_x^* + \mu H_z \cdot H_z^* \right\}$$

$$= \frac{1}{16\pi} \text{Re} \left\{ \epsilon E_i \cdot E_i \cdot e^{-i(ky+wt)} \cdot E_i \cdot e^{i(ky+wt)} + \mu \frac{\sqrt{\mu\epsilon} E_i \cdot E_i \cdot e^{-i(ky+wt)}}{\mu} \cdot \frac{\sqrt{\mu\epsilon} E_i \cdot E_i \cdot e^{i(ky+wt)}}{\mu} \right\}$$

$$Presión_{red} = \frac{1}{16\pi} \text{Re} \left\{ \epsilon E_i^2 + \epsilon E_i^2 \right\} = \frac{1}{16\pi} 2 E_i^2 \cdot \epsilon$$

* medio ②

$$d\vec{F} = \vec{T} \cdot d\vec{S} = \begin{pmatrix} T_{xx} & T_{yx} & T_{yz} \\ T_{yx} & T_{yy} & T_{yz} \\ T_{yz} & T_{yz} & T_{zz} \end{pmatrix} \begin{pmatrix} 0 \\ dx dz \\ 0 \end{pmatrix} = \begin{pmatrix} T_{yy} \cdot dx dz \end{pmatrix}$$

Presión = $\langle T_{yy} \rangle \Rightarrow$ análogamente a lo considerado para el medio 1 es:

$$T_{yy} = \frac{1}{4\pi} \left(-\frac{\epsilon' E_x^2}{2} - \frac{\mu' H_z^2}{2} \right) \therefore$$

$$\langle T_{yy} \rangle = \frac{1}{16\pi} \text{Re} \left\{ -\epsilon' E_x \cdot E_x^* - \mu' H_z \cdot H_z^* \right\}$$

$$= \frac{1}{16\pi} \text{Re} \left\{ -\epsilon' E_t^2 - \mu' \frac{\mu\epsilon}{\mu\epsilon'} E_t^2 \right\} = \frac{1}{16\pi} (-2\epsilon' E_t^2)$$

* Continuo

Bñ continuo

$$\sqrt{\mu\epsilon} E_i - \sqrt{\mu'\epsilon'} E_t = 0 \rightarrow E_t = E_i \left(\frac{\mu\epsilon}{\mu'\epsilon'} \right)^{1/2}$$

$$Presión = \frac{1}{16\pi} \epsilon 2 E_i^2 - \left(\frac{-1}{16\pi} \epsilon' 2 E_i^2 \cdot \frac{\mu\epsilon}{\mu'\epsilon'} \right) = \frac{1}{16\pi} E_i^2 \cdot 2 \cdot \epsilon \left(1 + \frac{\mu}{\mu'} \right)$$

$$\langle \mu_{em} \rangle = \frac{1}{16\pi} \operatorname{Re} \{ E \cdot D^* + B \cdot H^* \}$$

* medio 1

$$= \frac{1}{16\pi} \operatorname{Re} \left\{ \epsilon \cdot E_i^z + \frac{\mu \epsilon}{\mu'} E_i^z \right\}$$

* medio 2

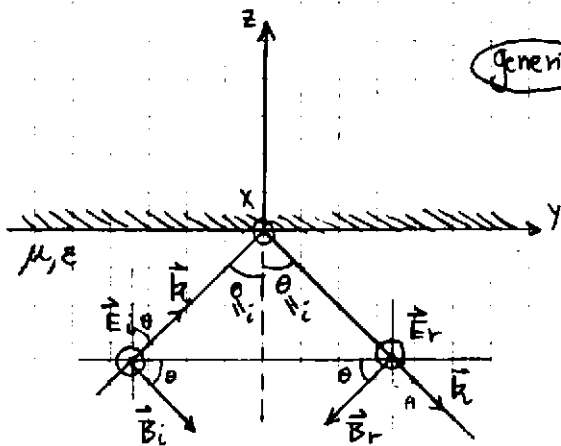
$$= \frac{1}{16\pi} \operatorname{Re} \left\{ \epsilon' E_i^z + \frac{\mu \epsilon'}{\mu'} E_i^z \right\}$$

$$\langle \mu_{em} \rangle = \frac{1}{16\pi} Z \epsilon E_i^z + \frac{1}{16\pi} Z \epsilon' \left(E_i^z \cdot \frac{\mu \epsilon}{\mu'} \right)$$

$$\boxed{\langle \mu_{em} \rangle = \frac{1}{16\pi} Z \epsilon E_i^z \left(1 + \frac{\mu}{\mu'} \right)}$$

▲ La densidad de energía es igual a la potencia de radiación

7. Análisis de las experiencias de Wiener



genérico a.

$$\begin{aligned}\vec{E}_i &= E_i \cdot e^{i(\vec{k} \cdot \vec{x} - \omega t)} \hat{x} \\ \vec{E}_r &= E_r \cdot e^{i(k_y y + k_z z - \omega t)} \hat{x} \\ \vec{B}_i &= \sqrt{\mu \epsilon} \cdot E_i \cdot e^{i(k_y y + k_z z - \omega t)} \\ &\quad (\cos \theta \hat{y} - \sin \theta \hat{z}) \\ \vec{E}_r &= E_r \cdot e^{i(k_y y - k_z z - \omega t)} \hat{x} \\ \vec{B}_r &= \sqrt{\mu \epsilon} \cdot E_r \cdot e^{i(k_y y - k_z z - \omega t)} \\ &\quad (-\cos \theta \hat{y} - \sin \theta \hat{z})\end{aligned}$$

* Contornos

Usamos $B_n = 0$ $E_z = 0 \rightarrow$

$$-\sqrt{\mu \epsilon} \cdot E_i \cdot \sin \theta - \sqrt{\mu \epsilon} \cdot E_r \cdot \sin \theta = 0 \Rightarrow E_i = -E_r$$

$$k_y = |\vec{k}| \cdot \sin \theta, \quad k_z = |\vec{k}| \cdot \cos \theta \quad (\text{K es la misma en cada cuadrante})$$

* Campos

$$\begin{aligned}\vec{E} &= \vec{E}_i + \vec{E}_r = E_i \cdot e^{i(k \sin \theta y - \omega t)} \begin{pmatrix} e^{i k \cos \theta z} & -e^{-i k \cos \theta z} \end{pmatrix} \hat{x} \\ \vec{B} &= \vec{B}_i + \vec{B}_r = \sqrt{\mu \epsilon} \cdot E_i \cdot e^{i(k \sin \theta y - \omega t)} \left[e^{i k \cos \theta z} (\cos \theta \hat{y} - \sin \theta \hat{z}) \right. \\ &\quad \left. - e^{-i k \cos \theta z} (-\cos \theta \hat{y} - \sin \theta \hat{z}) \right] \\ \vec{B} &= \sqrt{\mu \epsilon} \cdot E_i \cdot e^{i(k \sin \theta y - \omega t)} \begin{pmatrix} e^{i k \cos \theta z} + e^{-i k \cos \theta z} \end{pmatrix} \cos \theta \hat{y} \\ &\quad \sqrt{\mu \epsilon} \cdot E_i \cdot e^{i(k \sin \theta y - \omega t)} \begin{pmatrix} -e^{i k \cos \theta z} + e^{-i k \cos \theta z} \end{pmatrix} \sin \theta \hat{z}\end{aligned}$$

b.

$$\vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H}) = \frac{c}{4\pi \mu} (\vec{E} \times \vec{B})$$

$$\vec{E} = E_i \cdot e^{i(k \sin \theta y - \omega t)} \cdot 2i \cdot \sin(k \cos \theta z) \hat{x}$$

$$\begin{aligned}\vec{B} &= \sqrt{\mu \epsilon} \cdot E_i \cdot e^{i(k \sin \theta y - \omega t)} \cdot 2 \cdot \cos(k \cos \theta z) \cdot \cos \theta \hat{y} \\ &\quad - \sqrt{\mu \epsilon} \cdot E_i \cdot e^{i(k \sin \theta y - \omega t)} \cdot 2i \cdot \sin(k \cos \theta z) \cdot \sin \theta \hat{z}\end{aligned}$$

$$\vec{S} = \frac{c}{4\pi \mu} \cdot \left(2i E_i^2 \cdot e^{2i(k \sin \theta y - \omega t)} \cdot \sin(k \cos \theta z) \sqrt{\mu \epsilon} \cdot \cos(k \cos \theta z) \cdot \cos \theta \hat{z} \right. \\ \left. + 2i E_i^2 \cdot e^{2i(k \sin \theta y - \omega t)} \cdot \sin^2(k \cos \theta z) \sqrt{\mu \epsilon} \cdot \sin \theta \hat{y} \right)$$

$$\vec{S} = \frac{c}{4\pi\mu} \cdot 4E_i^2 \sqrt{\mu\epsilon} \cdot e^{i \cdot z \cdot (k \cos \theta - \omega t)} \left[\begin{array}{l} \sin(kz \cos \theta) \cos(kz \cos \theta) \cos \theta \cdot i \hat{z} \\ - \sin^2(kz \cos \theta) \sin \theta \hat{y} \end{array} \right]$$

$$\langle \vec{S} \rangle = \frac{c}{4\pi\mu} \cdot \frac{1}{z} \operatorname{Re}(\vec{E} \times \vec{B}^*) = \frac{c}{4\pi\mu} \cdot \frac{1}{z} \operatorname{Re}(\vec{E}^* \times \vec{B})$$

$$\vec{E}^* = E_i \cdot e^{-i(k \cos \theta y - \omega t + \pi/2)} \cdot 2 \sin(kz \cos \theta) \hat{x}$$

$$\vec{B} = \sqrt{\mu\epsilon} \cdot E_i \cdot 2 \cdot e^{+i(k \cos \theta y - \omega t)} \cdot \cos(k \cos \theta z) \cos \theta \hat{y} - \sqrt{\mu\epsilon} \cdot E_i \cdot 2 \cdot e^{i(k \cos \theta y - \omega t + \pi/2)} \cdot \sin(k \cos \theta z) \sin \theta \hat{z}$$

$$i = e^{i\pi/2}$$

$$\vec{E}^* \times \vec{B} = \begin{array}{l} 2^2 E_i^2 \cdot \sin(kz \cos \theta) \cdot \cos(k \cos \theta z) \cdot \cos \theta \cdot e^{-i\pi/2} \sqrt{\mu\epsilon} \hat{z} \\ 2^2 E_i^2 \cdot \sqrt{\mu\epsilon} \cdot \sin^2(kz \cos \theta) \cdot \sin \theta \hat{y} \end{array}$$

$$\langle \vec{S} \rangle = \frac{c}{8\pi\mu} \cdot 4E_i^2 \sqrt{\mu\epsilon} \sin^2(kz \cos \theta) \sin \theta \hat{y}$$

c.
$$\mu_e = \frac{1}{8\pi} (\vec{E} \cdot \vec{D})$$

$$\mu_e = \frac{1}{8\pi} \epsilon |\vec{E}|^2 = \frac{1}{8\pi} \epsilon 2^2 E_i^2 \sin^2(k \cos \theta z)$$

$$\langle \mu_e \rangle = \frac{1}{16\pi} \operatorname{Re}(\vec{E} \cdot \vec{D}^*) = \frac{1}{16\pi} \epsilon \cdot 4E_i^2 \sin^2(k \cos \theta z)$$

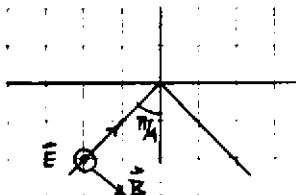
d.
$$\mu_m = \frac{1}{8\pi} (\vec{B} \cdot \vec{H})$$

$$\mu_m = \frac{1}{8\pi} \frac{1}{\mu} |\vec{B}|^2 = \frac{1}{8\pi\mu} \left(\mu\epsilon E_i^2 4 \cdot \cos^2(k \cos \theta z) \cdot \cos^2 \theta + \mu\epsilon E_i^2 4 \cdot \sin^2(k \cos \theta z) \cdot \sin^2 \theta \right)$$

$$\mu_m = \frac{1}{8\pi} \epsilon E_i^2 4 \left[\cos^2(k \cos \theta z) \cdot \cos^2 \theta + \sin^2(k \cos \theta z) \cdot \sin^2 \theta \right]$$

$$\langle \mu_m \rangle = \frac{1}{16\pi} \operatorname{Re}(\vec{B} \cdot \vec{H}^*) = \frac{1}{16\pi} \frac{1}{\mu} \epsilon E_i^2 4 \left[\cos^2(k \cos \theta z) \cos^2 \theta + \sin^2(k \cos \theta z) \sin^2 \theta \right]$$

②



Afortunadamente este es un caso particular del punto genérico y tan solo necesitamos especializar en $\theta = \frac{\pi}{4}$

$$\cos \pi/4 = \sin \pi/4 = a \rightarrow \sqrt{} \quad a = \frac{1}{\sqrt{2}}$$

a. * Campos

$$\vec{E} = E_i \cdot e^{i(kay - \omega t + \pi/2)} \cdot 2 \cdot \sin(kaz) \hat{x}$$

$$\vec{B} = \sqrt{\mu\epsilon} \cdot E_i \cdot e^{i(kay - \omega t)} \cdot 2 \cdot \cos(kaz) \cdot a \hat{y} - \sqrt{\mu\epsilon} \cdot E_i \cdot e^{i(kay - \omega t + \pi/2)} \cdot 2 \cdot \sin(kaz) \cdot a \hat{z}$$

b. * Vector de Poynting

$$\vec{S} = \frac{c}{4\pi\mu} \cdot 4E_i^2 \sqrt{\mu\epsilon} \cdot e^{i2(kya - \omega t)} \left(\sin(kaz) \cos(kaz) \cdot a \hat{z} - \sin^2(kaz) \cdot a \hat{y} \right)$$

$$\langle \vec{S} \rangle = \frac{c}{8\pi\mu} \cdot 4E_i^2 \sqrt{\mu\epsilon} \cdot a \cdot \sin^2(kaz) \hat{y}$$

c. * energía E

$$\mu_E = \frac{1}{8\pi} \epsilon 4E_i^2 \sin^2(kaz)$$

$$\langle \mu_E \rangle = \frac{1}{16\pi} \epsilon 4E_i^2 \sin^2(kaz)$$

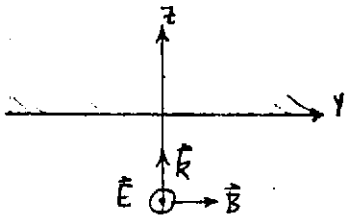
d. * energía M

$$\mu_M = \frac{1}{8\pi} \epsilon 4E_i^2 (\cos^2(kaz) \cdot a^2 + \sin^2(kaz) \cdot a^2)$$

$$\mu_M = \frac{1}{8\pi} \epsilon 4E_i^2 a^2$$

$$\langle \mu_M \rangle = \frac{1}{16\pi} \epsilon 4E_i^2 a^2$$

①



Incidencia normal es $\theta = 0 \Rightarrow$

$$\cos \theta = 1 \quad \sin \theta = 0$$

$$\vec{E} = E_i \cdot e^{-i(\omega t + \pi/2)} \cdot 2 \cdot \sin(kz) \hat{x}$$

$$\vec{B} = \sqrt{\mu\epsilon} \cdot E_i \cdot e^{-i(\omega t)} \cdot 2 \cdot \cos(kz) \hat{y}$$

$$\vec{S} = \frac{c}{4\pi\mu} \cdot 4E_i^2 \sqrt{\mu\epsilon} \cdot e^{-2i(\omega t + \pi/2)} \cdot \sin(kz) \cdot \cos(kz) \cdot \hat{z}$$

$$\langle \vec{S} \rangle = 0$$

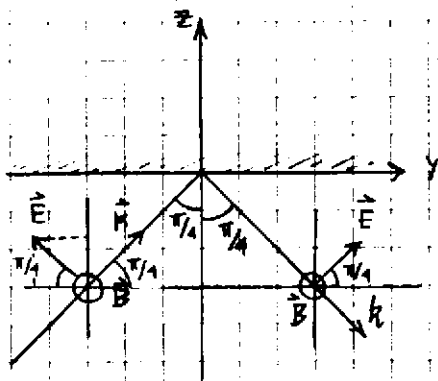
no hay flujo de energía en el tiempo

$$\langle \mu_E \rangle = \frac{1}{4\pi} \cdot \epsilon E_i^2 \cdot \sin^2(kz)$$

$$\langle \mu_M \rangle = \frac{1}{4\pi} \epsilon E_i^2 \cdot \cos^2(kz)$$

$$\langle \mu_{EM} \rangle = \frac{1}{4\pi} \epsilon E_i^2$$

3



$$\vec{E}_i = E_i e^{i(\vec{k} \cdot \vec{x} - \omega t)} a(\hat{z} - \hat{y})$$

$$\vec{E}_r = E_r e^{i(\vec{k} \cdot \vec{x} - \omega t)} a(\hat{z} + \hat{y})$$

$$\vec{B}_i = \sqrt{\mu\epsilon} E_i e^{i(\vec{k} \cdot \vec{x} - \omega t)} a \hat{x}$$

$$\vec{B}_r = \sqrt{\mu\epsilon} E_r e^{i(\vec{k} \cdot \vec{x} - \omega t)} a \hat{x}$$

* \perp

* \perp

k es el mismo para incidente y reflejado
* Lantornos

$$\vec{k} \cdot \vec{x} = k a y + k a z \quad \vec{k} \cdot \vec{x} = k a y - k a z$$

B_n, E_t continuos

$$-E_i a + E_r a = 0 \Rightarrow E_i = E_r$$

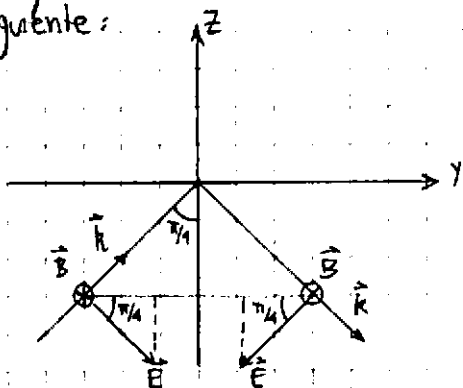
$$\vec{E} = \vec{E}_i + \vec{E}_r = E_i a e^{i(k a y - \omega t)} \left[e^{i k a z} (\hat{z} + \hat{y}) + e^{-i k a z} (\hat{z} - \hat{y}) \right]$$

$$E_i a e^{i(k a y - \omega t)} \left[(e^{i k a z} + e^{-i k a z}) \hat{z} \right. \\ \left. (-e^{i k a z} + e^{-i k a z}) \hat{y} \right]$$

$$\vec{E} = E_i a e^{i(k a y - \omega t)} 2 \left[(\cos(k a z)) \hat{z} + i \sin(k a z) \hat{y} \right]$$

$$\vec{B} = \vec{B}_i + \vec{B}_r = \sqrt{\mu\epsilon} E_i a e^{i(k a y - \omega t)} 2 \cos(k a z) \hat{x}$$

Podemos, llegados a este paso, entornos todo el trabajo de recalcular comparando con caso anterior, vemos que podríamos elegir ejes en el modo siguiente:



$$\vec{E}_i = E_i e^{i(k a y + k a z - \omega t)} a(-\hat{z} + \hat{y})$$

$$\vec{E}_r = E_r e^{i(k a y - k a z - \omega t)} a(\hat{z} - \hat{y})$$

$$\vec{B}_i = \sqrt{\mu\epsilon} E_i e^{i(k a y + k a z - \omega t)} a(-\hat{x})$$

$$\vec{B}_r = \sqrt{\mu\epsilon} E_r e^{i(k a y - k a z - \omega t)} a(\hat{x})$$

* Continuos

B_n, E_t continuos

a. * Campos

$$E_i a - E_r a = 0 \Rightarrow E_i = E_r$$

$$\vec{E} = \vec{E}_i + \vec{E}_r = E_i a e^{i(k a y - \omega t)} \left[-e^{i k a z} \hat{z} + e^{i k a z} \hat{y} - e^{-i k a z} \hat{z} - e^{-i k a z} \hat{y} \right]$$

$$\vec{E} = E_i a e^{i(k a y - \omega t)} \left[-\cos(k a z) 2 \hat{z} + 2i \sin(k a z) \hat{y} \right]$$

$$\vec{B} = \vec{B}_i + \vec{B}_r = -\sqrt{\mu\epsilon} \cdot E_i \cdot e^{i(kay - \omega t)} \cdot a [e^{ikaz} + e^{-ikaz}] \hat{x}$$

$$\boxed{\vec{B} = -\sqrt{\mu\epsilon} \cdot E_i \cdot a \cdot e^{i(kay - \omega t)} \cdot \cos(kaz) \cdot 2 \hat{x}}$$

Fue en vano; \vec{B}_{tm} es parecido en forma a \vec{E} del caso TE pero no igual y lo mismo con \vec{E}_{tm} que se parece a \vec{B}_{TE} . Hay que recalcular

b. * Vector de Poynting

$$\langle \vec{S} \rangle = \frac{c}{4\pi\mu} \frac{1}{z} \text{Re}(\vec{E} \times \vec{B}^*)$$

$$\vec{E} \times \vec{B}^* = -E_i^2 a^2 \sqrt{\mu\epsilon} \cdot 4 \cdot \cos^2(kaz) \hat{y} + E_i^2 a^2 \sqrt{\mu\epsilon} \cdot 4 \cdot e^{i\pi/2} \cdot \sin(kaz) \cdot \cos(kaz) \hat{z}$$

$$\boxed{\langle \vec{S} \rangle = \frac{c}{8\pi\mu} \cdot E_i^2 a^2 \cdot 4 \sqrt{\mu\epsilon} \cdot \cos^2(kaz) \hat{y}}$$

c., d. * Energía eléctrica y magnética

$$\langle \mu_E \rangle = \frac{1}{16\pi} \cdot \text{Re}(\epsilon |\vec{E}|^2) = \frac{\epsilon}{16\pi} \cdot (E_i^2 a^2 \cdot 4 [\cos^2(kaz) + \sin^2(kaz)])$$

$$\boxed{\langle \mu_E \rangle = \frac{\epsilon \cdot a^2 E_i^2 \cdot 4}{16\pi}}$$

$$\langle \mu_M \rangle = \frac{1}{16\pi} \cdot \text{Re}\left(\frac{1}{\mu} |\vec{B}|^2\right) = \frac{1}{16\pi} \frac{\mu\epsilon}{\mu} E_i^2 a^2 \cdot 4 \cdot \cos^2(kaz)$$

$$\boxed{\langle \mu_M \rangle = \frac{\epsilon}{16\pi} a^2 E_i^2 \cdot 4 \cdot \cos^2(kaz)}$$

* Conclusiones *

e. El vector óptico es \vec{E}

En la experiencia ③ no se observan franjas; es decir que la película no se impresionó con datos lumínicos. En dicha experiencia se tiene:

$\langle \mu_E \rangle$: constante

$\langle \mu_M \rangle$: depende de la coordenada z

En las experiencias ① y ② tenemos:

①

$\langle \mu_E \rangle$: depende de z

$\langle \mu_M \rangle$: depende de z

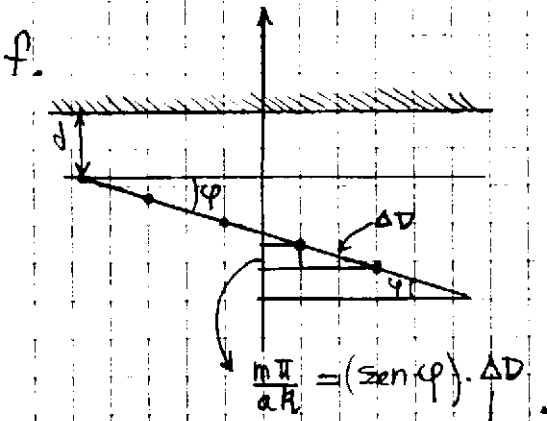
②

$\langle \mu_E \rangle$: depende de z

$\langle \mu_M \rangle$: constante

Luego como en ① y ② hubo franjas y en ③ no $\langle \mu_E \rangle$ es el causante de las franjas sobre la película; $\Rightarrow \vec{E}$ es el vector óptico

Asimismo en $z=0$ (pegado al espejo) $\langle \mu_M \rangle = 0$ lo cual sustenta aún más el argumento.

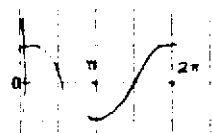


$$\frac{m\pi}{a.k} = (\text{sen } \varphi) \cdot \Delta D$$

$$\Delta D = \frac{1}{\text{sen } \varphi} \cdot \frac{m\pi}{a.k}$$

interferenja

$$\Delta D = \frac{1}{\text{sen } \varphi} \cdot m \cdot d$$



$$\cos^2(kaz) = 1 \Leftrightarrow 1$$

$$kaz = m\pi$$

$$z = \frac{m\pi}{a.k} = \frac{\sqrt{2} \cdot \pi \cdot m}{\sqrt{\mu \epsilon} \cdot \omega} \cdot c = \frac{m \cdot \lambda}{\sqrt{2}}$$

$$k = \frac{\omega}{c} \cdot \sqrt{\mu \epsilon}$$

$$z =$$

$$k = \frac{2\pi}{\lambda}$$

$$\lambda = \frac{c}{\nu}$$

$$z = -d \quad m=1$$

$$-d = -\frac{1}{a.k}$$